

Efficient Optimization of Aerodynamic Coefficients

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- **Aerodynamic forces and moments**
- **Aerodynamic shape optimization**
- **Compressible Euler equations**
- **Motivation of adjoint approach**
- **Continuous adjoint approach**
- **Derivation of adjoint Euler equations**
- **Implementation of an adjoint Euler solver**
- **Validation and application of continuous adjoint approach in 2D and 3D**

Aerodynamic Force

\vec{A}

\vec{L}

Lift

\vec{D}

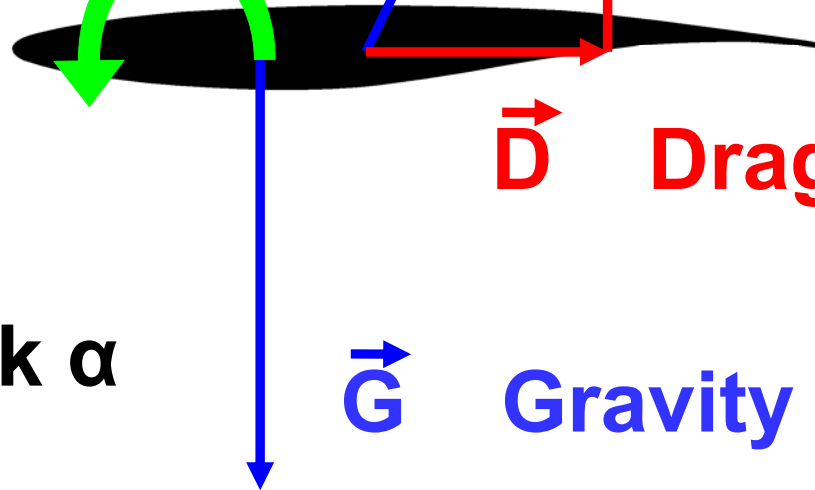
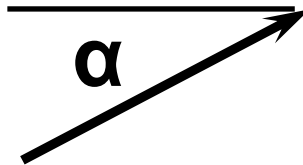
Drag

$\otimes \vec{M}$

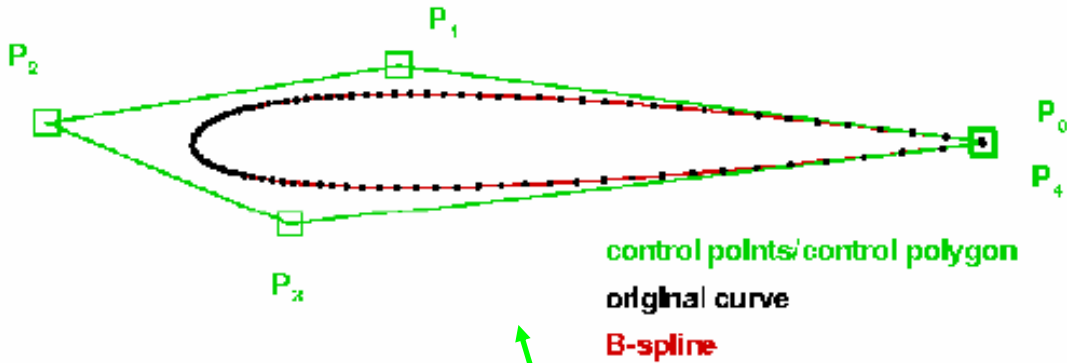
\vec{G}

Gravity

Pitching Moment



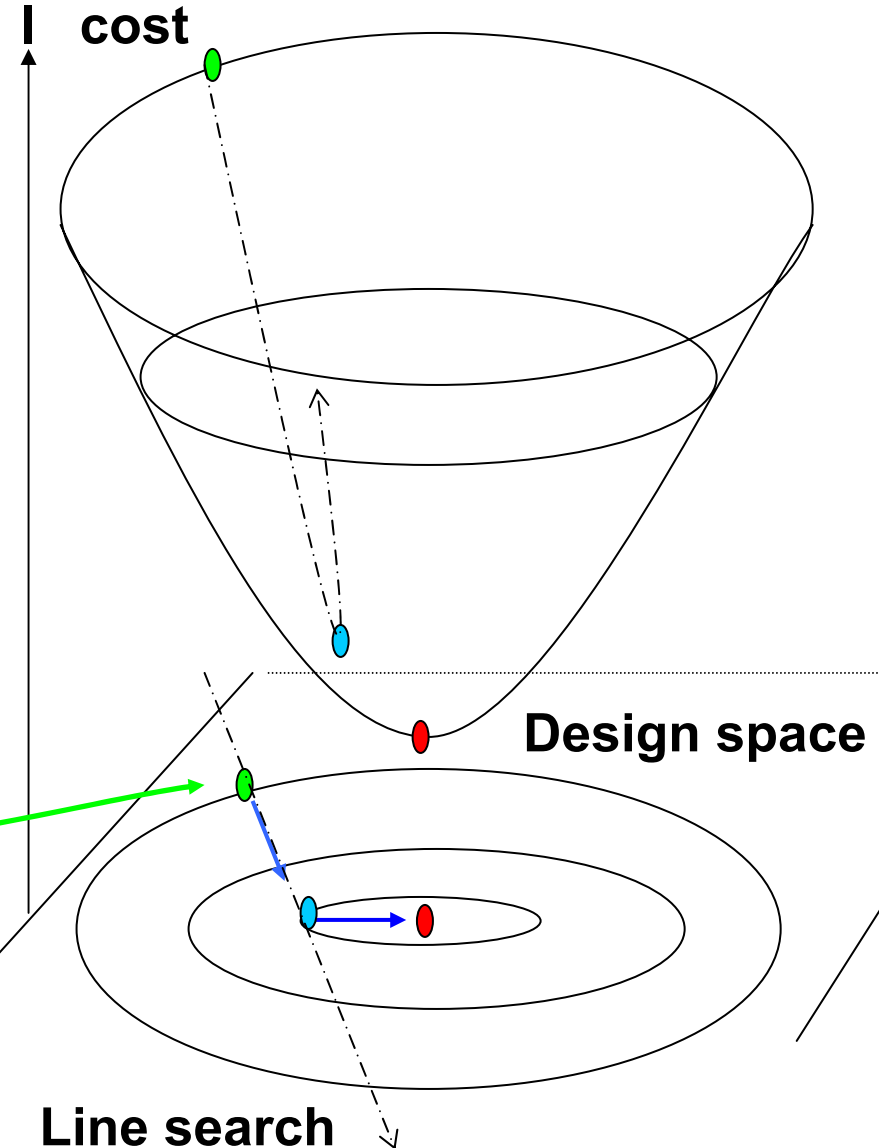
Angle of attack α



Parametrized airfoil

Search direction

$$- \nabla I = - \left(\dots, \frac{\delta I}{\delta P_i}, \dots \right)_{i=1, \dots, n}^T$$



Compressible 2D Euler Equations

$$\frac{\partial w}{\partial t} + \frac{\partial f}{\partial x} + \frac{\partial g}{\partial y} = 0$$

while

$$w = \begin{pmatrix} \rho \\ \rho u \\ \rho v \\ \rho E \end{pmatrix}, \quad f = \begin{pmatrix} \rho u \\ \rho u^2 + p \\ \rho uv \\ \rho uH \end{pmatrix}, \quad g = \begin{pmatrix} \rho v \\ \rho vu \\ \rho v^2 + p \\ \rho vH \end{pmatrix}$$

Pressure (ideal gas)

$$p = (\gamma - 1)\rho\left(E - \frac{1}{2}\vec{v}^2\right)$$

Dimensionless pressure

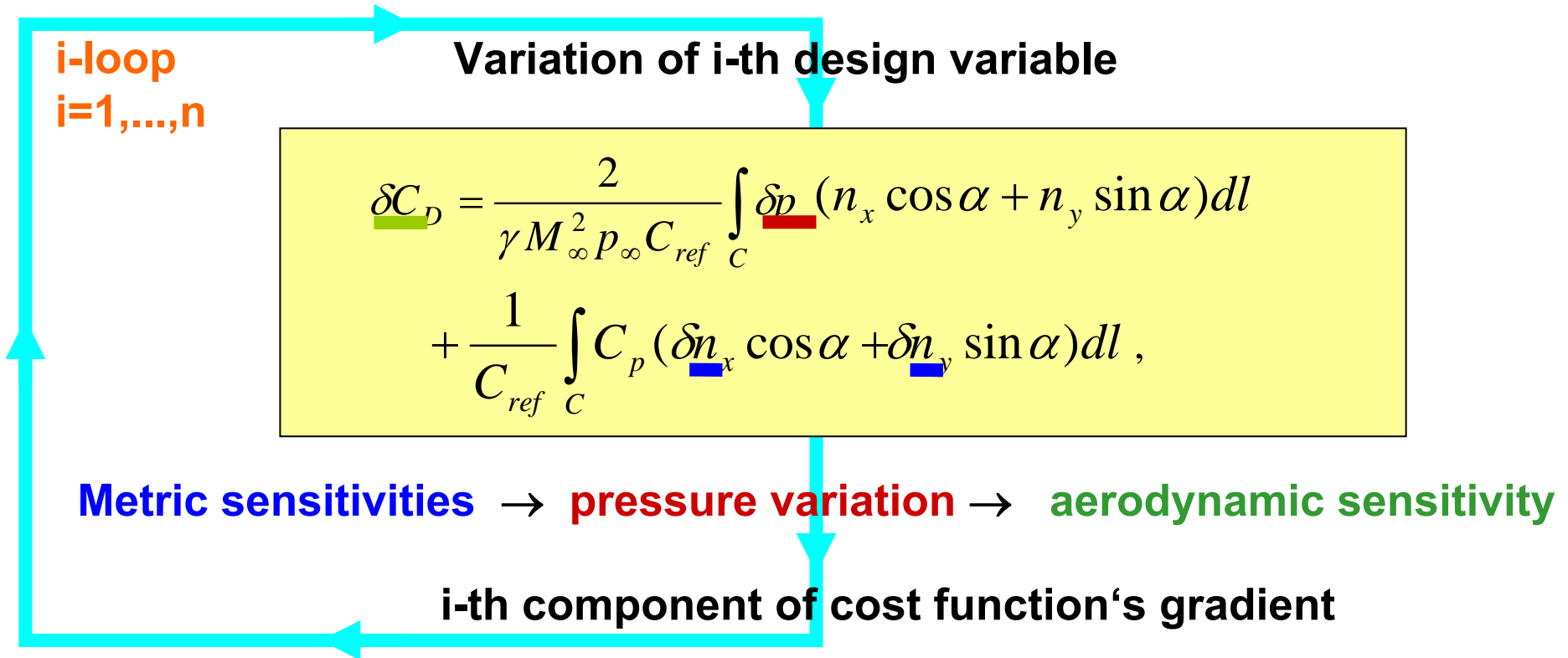
$$C_p = \frac{2(p - p_\infty)}{\gamma M_\infty^2 p_\infty}$$

Drag, lift, pitching moment coefficients

$$C_D = \frac{1}{C_{ref}} \int_C C_p (n_x \cos \alpha + n_y \sin \alpha) dl$$

$$C_L = \frac{1}{C_{ref}} \int_C C_p (n_y \cos \alpha - n_x \sin \alpha) dl$$

$$C_m = \frac{1}{C_{ref}^2} \int_C C_p (n_y (x - x_m) - n_x (y - y_m)) dl$$





• Finite Differences



n design variables require
n+1 flow calculations

High number of design variables

- **Finite Differences**  **n design variables require n+1 flow calculations**
- **Adjoint Approach**  **n design variables require 1 flow and 1 adjoint flow calculation**
Independent of number of design variables
High accuracy



Dual or Adjoint (Linear) Problem



Let be $A \in \mathbb{R}^{n \times m}$, $h \in \mathbb{R}^m$, $\varphi \in \mathbb{R}^m$ and $b \in \mathbb{R}^n$.

We define the primal linear problem:

$$\text{evaluate } I = h^T \varphi, \quad (1)$$

$$\text{while } A\varphi = b. \quad (2)$$

Furthermore, $\psi \in \mathbb{R}^n$ fulfills:

$$A^T \psi = h. \quad (3)$$

Then eqs. (2) and (3) imply

$$h^T \varphi = (A^T \psi)^T \varphi = (A^T \psi, \varphi) = (\psi, A\varphi) = \psi^T A\varphi = \psi^T b \quad \forall \varphi, \psi \quad (4)$$

and we have the equivalent dual or adjoint linear problem:

$$\text{evaluate } I = \psi^T b, \quad (5)$$

$$\text{while } A^T \psi = h. \quad (6)$$

The vector $\psi = (\psi_i)_{i \in \{1, \dots, n\}}$ is called the vector of adjoint variables ψ_i .

We define now the scalar product

$$(h, \varphi) := \int_{\Omega} h^T \varphi dx . \quad (7)$$

Let φ be the solution of the PDE

$$L\varphi = b \quad (8)$$

in the domain Ω , which fulfills the homogeneous boundary conditions on $\partial\Omega$.

Then L^* , the dual or adjoint operator of L , is defined as:

$$L^* : (\psi, L\varphi) = (L^*\psi, \varphi) \quad \forall \varphi, \psi. \quad (9)$$

Furthermore, ψ , the vector(-field) of adjoint variables, solves the dual or adjoint PDE

$$L^*\psi = h \quad (10)$$

in the domain Ω and again fulfills the homogeneous boundary conditions on $\partial\Omega$.

Then finally we have as before:

$$(h, \varphi) = (L^*\psi, \varphi) = (\psi, L\varphi) = (\psi, b). \quad (11)$$

Let's take e.g. the convection-diffusion equation

$$L\varphi \equiv \frac{d\varphi}{dx} - \epsilon \frac{d^2\varphi}{dx^2}, \quad 0 < x < 1, \quad (12)$$

with homogeneous boundary conditions $\varphi(0) = \varphi(1) = 0$.

Integration by parts yields ($\varphi, \psi \in C^2$):

$$(\psi, L\varphi) = \int_0^1 \psi \left(\frac{d\varphi}{dx} - \epsilon \frac{d^2\varphi}{dx^2} \right) dx \quad (13)$$

$$= \int_0^1 \left(-\frac{d\psi}{dx} - \epsilon \frac{d^2\psi}{dx^2} \right) \varphi dx + \left[\psi\varphi - \epsilon\psi \frac{d\varphi}{dx} + \epsilon\varphi \frac{d\psi}{dx} \right]_0^1 \quad (14)$$

$$= \int_0^1 \underbrace{\left(-\frac{d\psi}{dx} - \epsilon \frac{d^2\psi}{dx^2} \right)}_{=: L^*\psi} \varphi dx + \left[-\epsilon\psi \frac{d\varphi}{dx} \right]_0^1. \quad (15)$$

For the adjoint convection-diffusion equation

$$L^*\psi \equiv -\frac{d\psi}{dx} - \epsilon \frac{d^2\psi}{dx^2}, \quad (16)$$

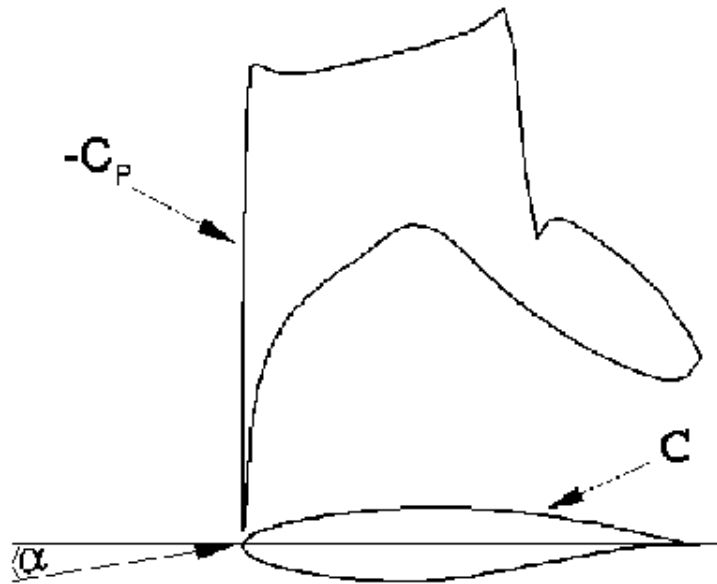
with homogeneous boundary conditions $\psi(0) = \psi(1) = 0$, the boundary term (15) vanishes and it holds (11):

$$(h, \varphi) = (L^*\psi, \varphi) = (\psi, L\varphi) = (\psi, b).$$

Some examples:

	Operator	Adjoint
Convection-Diffusion Eq.	$\frac{d\varphi}{dx} - \epsilon \frac{d^2\varphi}{dx^2}$	$-\frac{d\psi}{dx} - \epsilon \frac{d^2\psi}{dx^2}$
Wave Eq.	$\frac{d\varphi}{dt} - \frac{d^2\varphi}{dx^2}$	$-\frac{d\psi}{dt} - \frac{d^2\psi}{dx^2}$
Convection Eq.	$\frac{d\varphi}{dt} + \frac{d\varphi}{dx}$	$-\frac{d\psi}{dt} - \frac{d\psi}{dx}$

Derivation of the Adjoint Euler Equations



D flow field domain

B far field

C wall

$\partial D := B \cup C$ flow field boundary

$\vec{S} := \begin{pmatrix} S_x \\ S_y \end{pmatrix}$ normal vector $\perp \partial D$

$\vec{n} := \begin{pmatrix} n_1 \\ n_2 \end{pmatrix}$ normal unit vector $\perp \partial D$

α angle of attack

C_D drag coefficient

C_L lift coefficient

p pressure

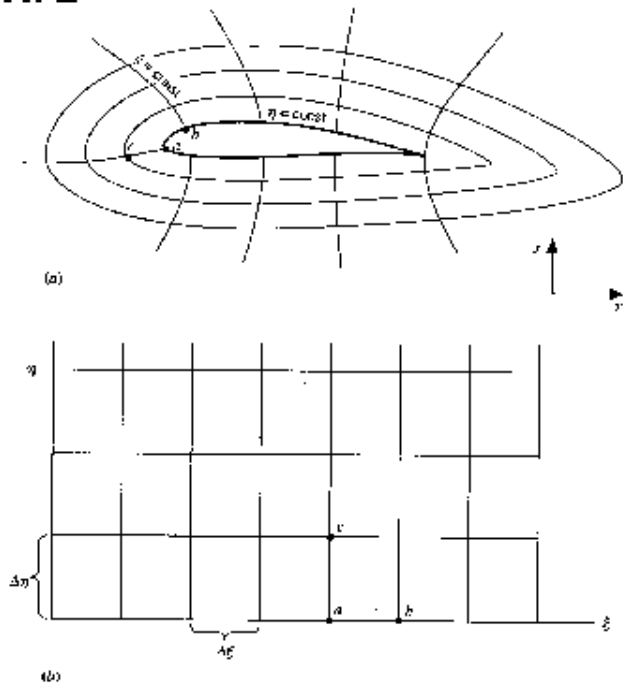
M Mach number

\dots_∞ ... at free stream

γ ratio of specific heats

S_{ref} area of airfoil

$\frac{2(p-p_\infty)}{\gamma M_\infty^2 p_\infty} =: C_p$ pressure coefficient



Cartesian coordinates:

$$\frac{\partial w}{\partial t} + \frac{\partial f}{\partial x} + \frac{\partial g}{\partial y} = 0$$

$$w = \begin{pmatrix} \rho \\ \rho u \\ \rho v \\ \rho E \end{pmatrix}, f = \begin{pmatrix} \rho u \\ \rho u^2 + p \\ \rho uv \\ \rho u H \end{pmatrix}, g = \begin{pmatrix} \rho v \\ \rho v u \\ \rho v^2 + p \\ \rho v H \end{pmatrix}$$

$$p = (\gamma - 1)\rho(E - \frac{1}{2}(u^2 + v^2)), \quad \rho H = \rho E + p$$

Body fitted transformation:

$$(x, y) \mapsto (\xi(x, y), \eta(x, y)),$$

$$J = \det \begin{pmatrix} \frac{\partial x}{\partial \xi} & \frac{\partial x}{\partial \eta} \\ \frac{\partial y}{\partial \xi} & \frac{\partial y}{\partial \eta} \end{pmatrix}, \quad \begin{pmatrix} U \\ V \end{pmatrix} = \frac{1}{J} \begin{pmatrix} \frac{\partial y}{\partial \eta} & -\frac{\partial x}{\partial \eta} \\ -\frac{\partial y}{\partial \xi} & \frac{\partial x}{\partial \xi} \end{pmatrix} \begin{pmatrix} u \\ v \end{pmatrix}$$

Body fitted coordinates:

$$\frac{\partial W}{\partial t} + \frac{\partial F}{\partial \xi} + \frac{\partial G}{\partial \eta} = 0$$

$$W = J \begin{pmatrix} \rho \\ \rho u \\ \rho v \\ \rho E \end{pmatrix}, F = J \begin{pmatrix} \rho U \\ \rho U u + \frac{\partial \xi}{\partial x} p \\ \rho U v + \frac{\partial \xi}{\partial y} p \\ \rho U H \end{pmatrix}, G = J \begin{pmatrix} \rho V \\ \rho V u + \frac{\partial \eta}{\partial x} p \\ \rho V v + \frac{\partial \eta}{\partial y} p \\ \rho V H \end{pmatrix}$$

In the case of steady state it holds for the perturbed geometry

$$\frac{\partial}{\partial \xi}(F + \delta F) + \frac{\partial}{\partial \eta}(G + \delta G) = 0$$
$$\Rightarrow (1) \quad \frac{\partial}{\partial \xi}(\delta F) + \frac{\partial}{\partial \eta}(\delta G) \stackrel{!}{=} 0.$$

Furthermore

$$(2) \quad \delta F = \delta \left(J \frac{\partial \xi}{\partial x} \right) f + \delta \left(J \frac{\partial \xi}{\partial y} \right) g + J \frac{\partial \xi}{\partial x} \frac{\partial f}{\partial w} \delta w + J \frac{\partial \xi}{\partial y} \frac{\partial g}{\partial w} \delta w$$

and

$$(3) \quad \delta G = \delta \left(J \frac{\partial \eta}{\partial x} \right) f + \delta \left(J \frac{\partial \eta}{\partial y} \right) g + J \frac{\partial \eta}{\partial x} \frac{\partial f}{\partial w} \delta w + J \frac{\partial \eta}{\partial y} \frac{\partial g}{\partial w} \delta w.$$

Together with (1) and the fundamental lemma of variational calculus it holds

$$\int_D \psi^T \left(\frac{\partial}{\partial \xi} (\delta F) + \frac{\partial}{\partial \eta} (\delta G) \right) d\xi d\eta = 0$$

for any Lagrangian multiplier ψ .

If ψ is differentiable one obtains together with Greens formula

$$-\int_D \left(\frac{\partial \psi^T}{\partial \xi} \delta F + \frac{\partial \psi^T}{\partial \eta} \delta G \right) d\xi d\eta + \int_B (n_1 \psi^T \delta F + n_2 \psi^T \delta G) d\xi - \int_C (n_1 \psi^T \delta F + n_2 \psi^T \delta G) d\xi = 0.$$

Now the variation of the cost function can be expressed as

$$\begin{aligned} \delta C_D &= \frac{2}{\gamma M_\infty^2 p_\infty S_{ref}} \int_C \delta p (S_x \cos \alpha + S_y \sin \alpha) d\xi - \int_D \left(\frac{\partial \psi^T}{\partial \xi} \delta F + \frac{\partial \psi^T}{\partial \eta} \delta G \right) d\xi d\eta \\ &+ \int_B (n_1 \psi^T \delta F + n_2 \psi^T \delta G) d\xi - \int_C \underbrace{(n_1 \psi^T \delta F + n_2 \psi^T \delta G)}_{=0, n_1=0} d\xi \\ &+ \frac{1}{S_{ref}} \int_C C_p (\delta S_x \cos \alpha + \delta S_y \sin \alpha) d\xi. \end{aligned}$$

Along C it holds $V = 0$ and yields

$$G = J \begin{pmatrix} 0 \\ \frac{\partial \eta}{\partial x} p \\ \frac{\partial \eta}{\partial y} p \\ 0 \end{pmatrix}, \quad \delta G = J \begin{pmatrix} 0 \\ \frac{\partial \eta}{\partial x} \delta p \\ \frac{\partial \eta}{\partial y} \delta p \\ 0 \end{pmatrix} + p \begin{pmatrix} 0 \\ \delta \left(J \frac{\partial \eta}{\partial x} \right) \\ \delta \left(J \frac{\partial \eta}{\partial y} \right) \\ 0 \end{pmatrix}.$$

Together with (2) and (3) one obtains

$$\begin{aligned}
 \delta C_D = & \frac{2}{\gamma M_\infty^2 p_\infty S_{ref}} \int_C \delta p (S_x \cos \alpha + S_y \sin \alpha) d\xi \\
 & - \int_D \frac{\partial \psi^T}{\partial \xi} \left(\delta \left(J \frac{\partial \xi}{\partial x} \right) f + \delta \left(J \frac{\partial \xi}{\partial y} \right) g + J \frac{\partial \xi}{\partial x} \frac{\partial f}{\partial w} \delta w + J \frac{\partial \xi}{\partial y} \frac{\partial g}{\partial w} \delta w \right) \\
 & \quad + \frac{\partial \psi^T}{\partial \eta} \left(\delta \left(J \frac{\partial \eta}{\partial x} \right) f + \delta \left(J \frac{\partial \eta}{\partial y} \right) g + J \frac{\partial \eta}{\partial x} \frac{\partial f}{\partial w} \delta w + J \frac{\partial \eta}{\partial y} \frac{\partial g}{\partial w} \delta w \right) d\xi d\eta \\
 & - \int_C \psi_2 \left(J \frac{\partial \eta}{\partial x} \delta p + p \delta \left(J \frac{\partial \eta}{\partial x} \right) \right) + \psi_3 \left(J \frac{\partial \eta}{\partial y} \delta p + p \delta \left(J \frac{\partial \eta}{\partial y} \right) \right) d\xi \\
 & + \int_B n_1 \psi^T \delta F + n_2 \psi^T \delta G d\xi + \frac{1}{S_{ref}} \int_C C_p (\delta S_x \cos \alpha + \delta S_y \sin \alpha) d\xi.
 \end{aligned}$$

If the adjoint Euler equations

$$\frac{\partial \psi^T}{\partial \xi} \left(J \frac{\partial \xi}{\partial x} \frac{\partial f}{\partial w} + J \frac{\partial \xi}{\partial y} \frac{\partial g}{\partial w} \right) + \frac{\partial \psi^T}{\partial \eta} \left(J \frac{\partial \eta}{\partial x} \frac{\partial f}{\partial w} + J \frac{\partial \eta}{\partial y} \frac{\partial g}{\partial w} \right) = 0 \quad \Leftrightarrow \quad \left(\frac{\partial f}{\partial w} \right)^T \frac{\partial \psi}{\partial x} + \left(\frac{\partial g}{\partial w} \right)^T \frac{\partial \psi}{\partial y} = 0$$

... are fulfilled in the domain D with the boundary conditions

$$\frac{2}{\gamma M_{\infty}^2 p_{\infty} S_{ref}} (S_x \cos \alpha + S_y \sin \alpha) = \underbrace{-S_x \psi_2 - S_y \psi_3}_{-\frac{\partial y}{\partial \xi} \psi_2 + \frac{\partial x}{\partial \xi} \psi_3 = J \frac{\partial \eta}{\partial x} \psi_2 + J \frac{\partial \eta}{\partial y} \psi_3}$$

on the airfoil C (dependent on the cost function!) and

$$\delta \left(J \frac{\partial \xi}{\partial x} \right), \dots, \delta \left(J \frac{\partial \eta}{\partial y} \right) \rightarrow 0 \quad \psi^T J \frac{\partial \xi}{\partial x} \frac{\partial f}{\partial w} \delta w = 0, \dots, \psi^T J \frac{\partial \eta}{\partial y} \frac{\partial g}{\partial w} \delta w = 0$$

at the far field B one can simplify δC_D to

$$\begin{aligned} \delta C_D = & - \int_D \frac{\partial \psi^T}{\partial \xi} \left(\delta \left(\frac{\partial y}{\partial \eta} \right) f - \delta \left(\frac{\partial x}{\partial \eta} \right) g \right) + \frac{\partial \psi^T}{\partial \eta} \left(-\delta \left(\frac{\partial y}{\partial \xi} \right) f + \delta \left(\frac{\partial x}{\partial \xi} \right) g \right) d\xi d\eta \\ & - \int_C p (\delta S_x \psi_2 + \delta S_y \psi_3) d\xi + \frac{1}{S_{ref}} \int_C C_p (\delta S_x \cos \alpha + \delta S_y \sin \alpha) d\xi. \end{aligned}$$

**Adjoint
Euler-Equations:**

$$-\frac{\partial \psi}{\partial t} - \left(\frac{\partial f}{\partial w} \right)^T \frac{\partial \psi}{\partial x} - \left(\frac{\partial g}{\partial w} \right)^T \frac{\partial \psi}{\partial y} = 0$$

Ψ : Vector of adjoint variables

Boundary conditions:

Wall: $n_x \psi_2 + n_y \psi_3 = \underline{-d(I)}$

Farfield: $\delta x_\xi, \dots, \delta y_\eta = 0, \delta w = 0$

Adjoint formulation of cost function's gradient:

$$\delta I = - \int_C p (-\psi_2 \delta y_\xi + \psi_3 \delta x_\xi) dl + \underline{K(I)}$$

$$- \int_D \psi_\xi^T (\delta y_\eta f - \delta x_\eta g) + \psi_\eta^T (-\delta y_\xi f + \delta x_\xi g) dA$$

$$d(C_D) = \frac{2}{\gamma M_\infty^2 p_\infty C_{ref}} (n_x \cos \alpha + n_y \sin \alpha)$$

Drag

$$K(C_D) = \frac{1}{C_{ref}} \int_C C_p (\delta n_x \cos \alpha + \delta n_y \sin \alpha) dl$$

$$d(C_L) = \frac{2}{\gamma M_\infty^2 p_\infty C_{ref}} (n_y \cos \alpha - n_x \sin \alpha)$$

Lift

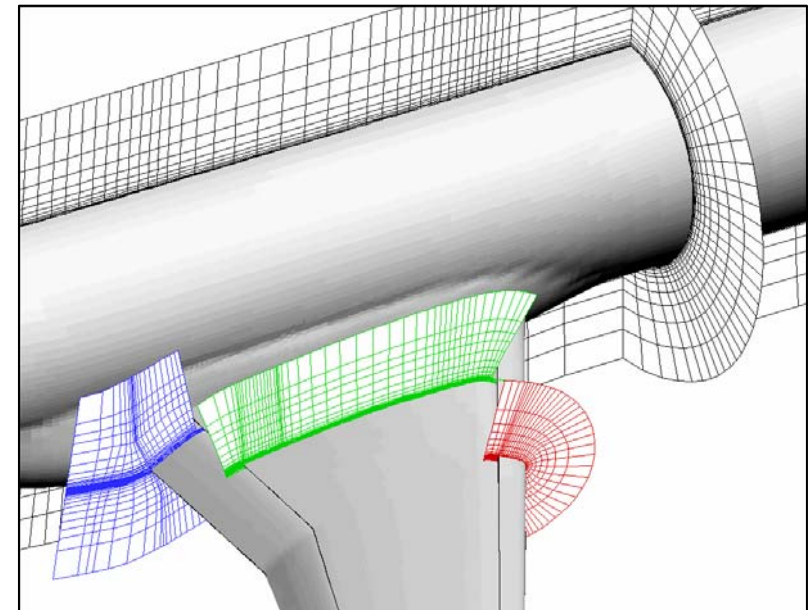
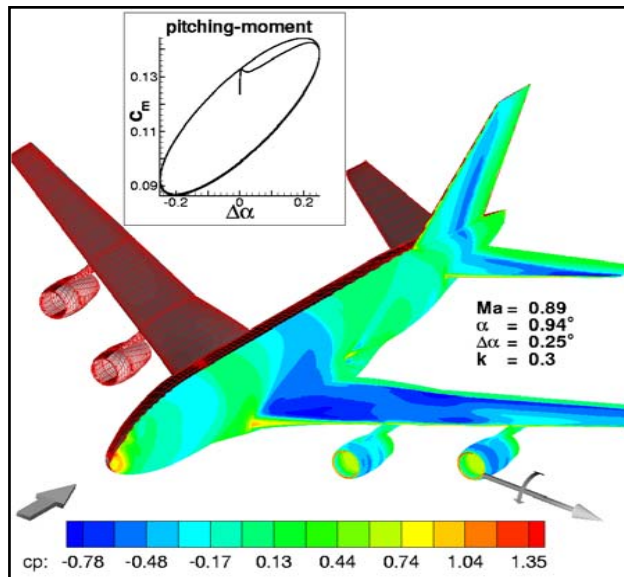
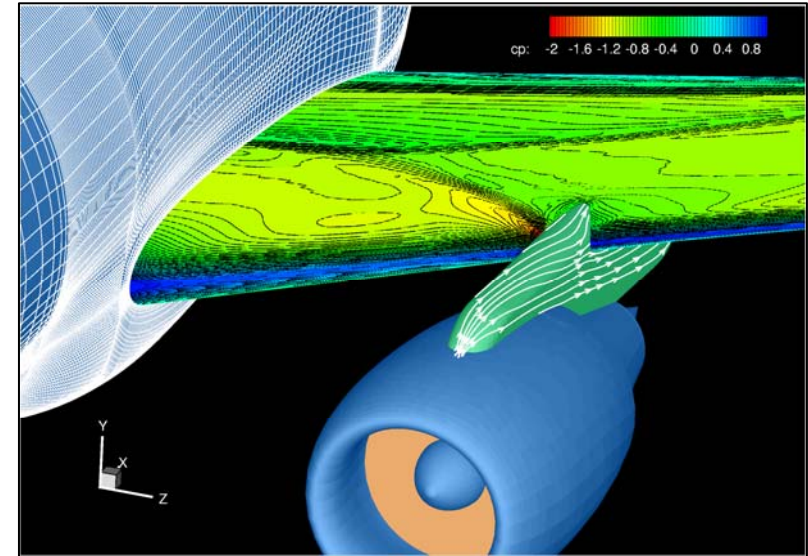
$$K(C_L) = \frac{1}{C_{ref}} \int_C C_p (\delta n_y \cos \alpha - \delta n_x \sin \alpha) dl$$

$$d(C_m) = \frac{2}{\gamma M_\infty^2 p_\infty C_{ref}^2} (n_y (x - x_m) - n_x (y - y_m))$$

Pitching moment

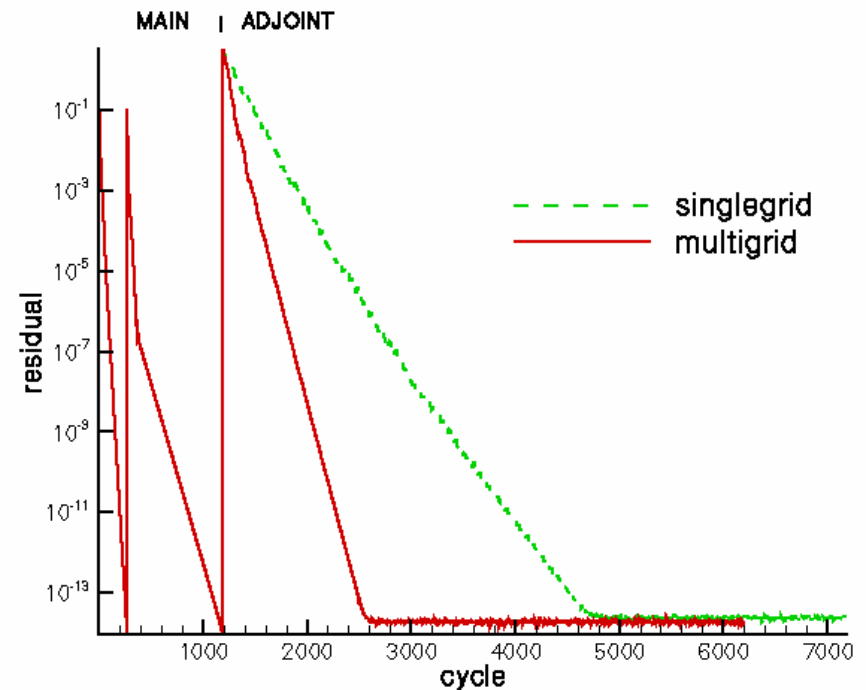
$$K(C_m) = \frac{1}{C_{ref}^2} \int_C C_p \delta (n_y (x - x_m) - n_x (y - y_m)) dl$$

- advanced turbulence and transition models
- 2nd order finite volume discretization (cell centered & cell vertex option)
- central and upwind schemes
- multigrid
- implicit treatment of turbulence equations
- implicit schemes for time accurate flows
- preconditioning for low speed flow
- design option (inverse design, **adjoint**)

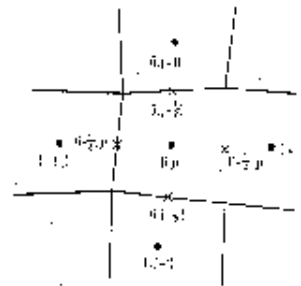
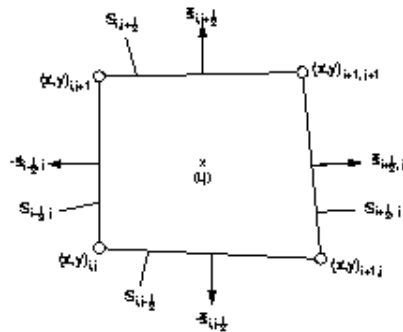


Adjoint solver on block-structured grids

- continuous adjoint approach
- implemented in **FLOWer**
- cost functions: lift, drag & moment and combinations
- adjoint solver based on multigrid
- Euler & Navier-Stokes (frozen μ)
- AD for turbulence equations (*FastOpt*)



convergence history, FLOWer



$$\frac{\partial}{\partial t} \int_{V_{i,j}} \psi dV + \int_{\partial V_{i,j}} \bar{F} \cdot \vec{n} ds = 0, \quad \bar{F} = \left[\left(\frac{\partial f}{\partial w} \right)^T \psi, \left(\frac{\partial g}{\partial w} \right)^T \psi \right]$$

$$\psi_{i,j} := \frac{1}{V_{i,j}} \int_{V_{i,j}} \psi dV, \quad \text{e.g. } \bar{F}_{i,j+\frac{1}{2}} := \left[\left(\frac{\partial f}{\partial w} \right)_{i,j}^T \frac{\psi_{i,j+1} + \psi_{i,j}}{2}, \left(\frac{\partial g}{\partial w} \right)_{i,j}^T \frac{\psi_{i,j+1} + \psi_{i,j}}{2} \right],$$

$$\text{Adjoint flux: } \bar{Q}_{i,j} := \bar{F}_{i,j+\frac{1}{2}} \bar{S}_{i,j+\frac{1}{2}} - \bar{F}_{i,j-\frac{1}{2}} \bar{S}_{i,j-\frac{1}{2}} + \bar{F}_{i+\frac{1}{2},j} \bar{S}_{i+\frac{1}{2},j} - \bar{F}_{i-\frac{1}{2},j} \bar{S}_{i-\frac{1}{2},j},$$

$$\left(\begin{array}{c} \bar{s}_x^{(j)} \\ \bar{s}_y^{(j)} \end{array} \right) := \frac{\bar{S}_{i,j+\frac{1}{2}} + \bar{S}_{i,j-\frac{1}{2}}}{2}, \quad \left(\begin{array}{c} \bar{s}_x^{(i)} \\ \bar{s}_y^{(i)} \end{array} \right) := \frac{\bar{S}_{i+\frac{1}{2},j} + \bar{S}_{i-\frac{1}{2},j}}{2},$$

$$\Rightarrow Q_{i,j} = \left(\bar{s}_x^{(j)} \left(\frac{\partial f}{\partial w} \right)_{i,j}^T + \bar{s}_y^{(j)} \left(\frac{\partial g}{\partial w} \right)_{i,j}^T \right) \frac{\psi_{i,j+1} - \psi_{i,j-1}}{2} + \left(\bar{s}_x^{(i)} \left(\frac{\partial f}{\partial w} \right)_{i,j}^T + \bar{s}_y^{(i)} \left(\frac{\partial g}{\partial w} \right)_{i,j}^T \right) \frac{\psi_{i+1,j} - \psi_{i-1,j}}{2}$$

$$V_{i,j} \frac{\partial \psi_{i,j}}{\partial t} + Q_{i,j} + D_{i,j} = 0, \quad D_{i,j}: \text{artificial dissipation}$$

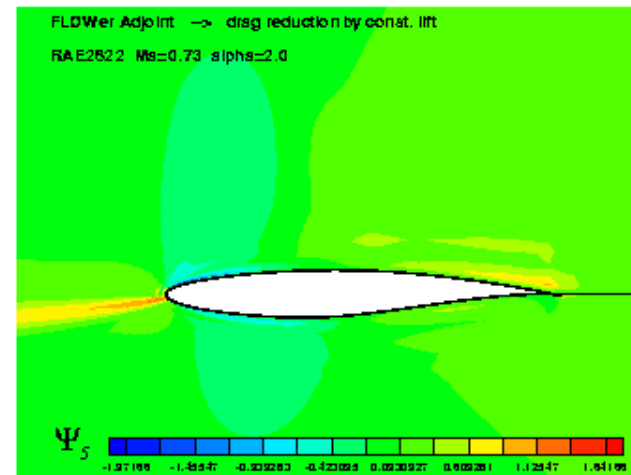
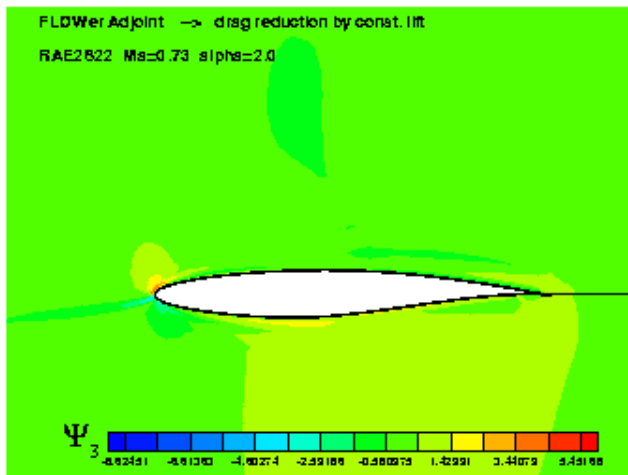
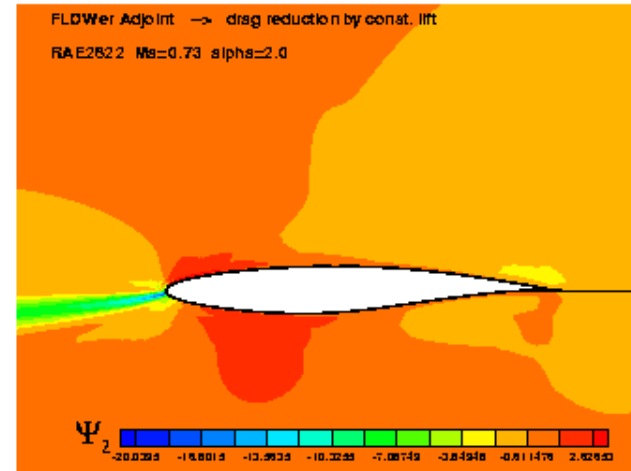
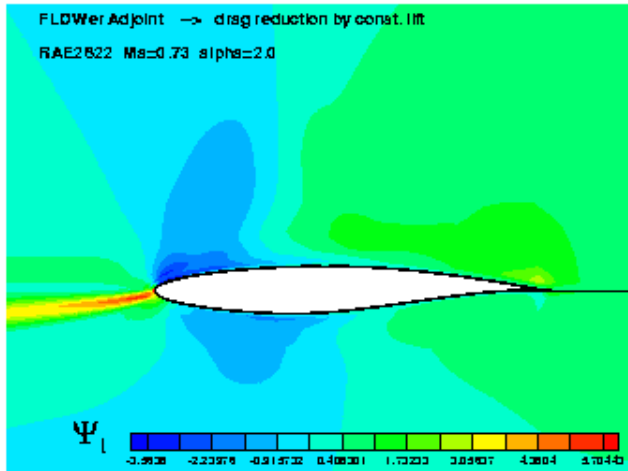
$$Q_{i,j} = \left(\bar{s}_x^{(j)} \left(\frac{\partial f}{\partial w} \right)_{i,j}^T + \bar{s}_y^{(j)} \left(\frac{\partial g}{\partial w} \right)_{i,j}^T \right) \frac{\psi_{i,j+1} - \psi_{i,j-1}}{2} + \left(\bar{s}_x^{(i)} \left(\frac{\partial f}{\partial w} \right)_{i,j}^T + \bar{s}_y^{(i)} \left(\frac{\partial g}{\partial w} \right)_{i,j}^T \right) \frac{\psi_{i+1,j} - \psi_{i-1,j}}{2}$$

$$\underbrace{T^{-1} T \left(\bar{s}_x^{(j)} \left(\frac{\partial f}{\partial w} \right)_{i,j}^T + \bar{s}_y^{(j)} \left(\frac{\partial g}{\partial w} \right)_{i,j}^T \right) T^{-1} T}_{=: Q_{i,j}^{(j)}} \hat{\psi}_{i,j}^{(j)} := T \frac{\psi_{i,j+1} - \psi_{i,j-1}}{2},$$

$$Q_{i,j}^{(j)} = \begin{pmatrix} u\bar{s}_x^{(j)} + v\bar{s}_y^{(j)} & 0 & 0 & 0 \\ \bar{s}_x^{(j)} & u\bar{s}_x^{(j)} + v\bar{s}_y^{(j)} & 0 & \frac{\gamma}{\gamma-1} \frac{p}{\rho} \bar{s}_x^{(j)} \\ \bar{s}_y^{(j)} & 0 & u\bar{s}_x^{(j)} + v\bar{s}_y^{(j)} & \frac{\gamma}{\gamma-1} \frac{p}{\rho} \bar{s}_y^{(j)} \\ 0 & (\gamma-1)\bar{s}_x^{(j)} & (\gamma-1)\bar{s}_y^{(j)} & u\bar{s}_x^{(j)} + v\bar{s}_y^{(j)} \end{pmatrix},$$

$$T^{-1} = \begin{pmatrix} 1 & -u & -v & \frac{1}{2}(u^2 + v^2) \\ 0 & 1 & 0 & -u \\ 0 & 0 & 1 & -v \\ 0 & 0 & 0 & 1 \end{pmatrix}, \quad \Rightarrow Q_{i,j} = T^{-1} \left(Q_{i,j}^{(j)} \hat{\psi}_{i,j}^{(j)} + Q_{i,j}^{(i)} \hat{\psi}_{i,j}^{(i)} \right)$$

Adjoint vector calculated by FLOWer



'Grid moving technique' based on J. Reuther

$$\delta C_D = - \int_D \frac{\partial \psi^T}{\partial \xi} (\delta y_{\eta} f - \delta x_{\eta} g) + \frac{\partial \psi^T}{\partial \eta} (-\delta y_{\xi} f + \delta x_{\xi} g) d\xi d\eta$$

$$- \int_C p (-\psi_2 \delta y_{\xi_s} + \psi_3 \delta x_{\xi_s}) d\xi + \frac{1}{S_{ref}} \int_C C_p (\delta y_{\xi_s} \cos \alpha - \delta x_{\xi_s} \sin \alpha) d\xi$$

Distance function e.g. $R(\eta) = 1 - \frac{\eta}{\eta_B}$

Values at wall C : $\eta = 0 \rightarrow R(0) = 1$

Values at far field B : $\eta = \eta_B \rightarrow R(\eta_B) = 0$

The grid for the perturbed (new) geometry is defined as:

$$x_{new} := x_{old} + R \cdot (x_{new_s} - x_{old_s}),$$

$$y_{new} := y_{old} + R \cdot (y_{new_s} - y_{old_s}).$$

Index s means the values at the surface.

Now the metric sensitivities can be expressed as

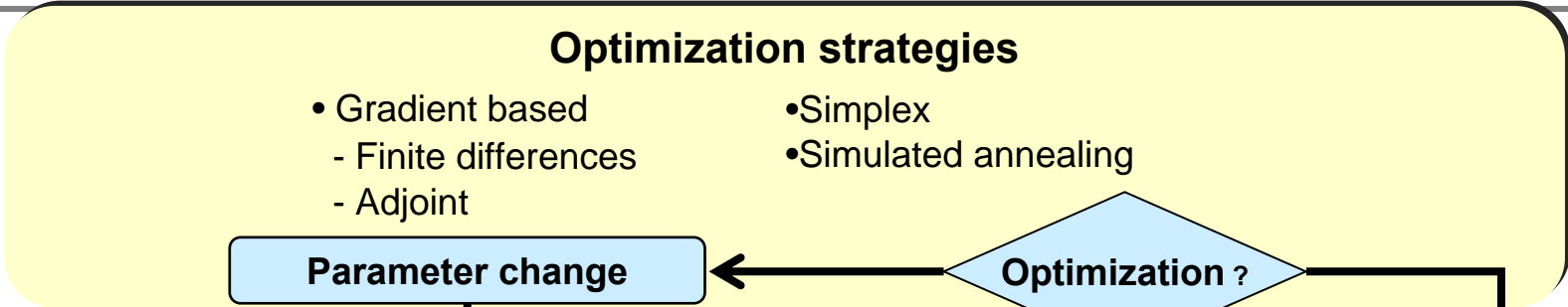
$$\delta x_{\xi} \approx R \cdot \delta x_{\xi_s}, \quad \delta y_{\xi} \approx R \cdot \delta y_{\xi_s},$$

$$\delta x_{\eta} \approx R \cdot \delta x_{\eta_s}, \quad \delta y_{\eta} \approx R \cdot \delta y_{\eta_s}$$

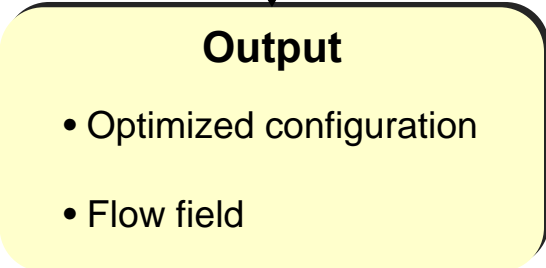
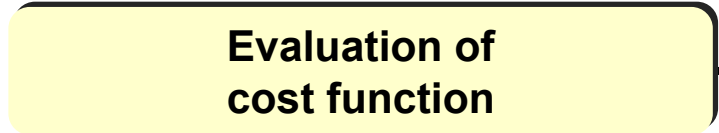
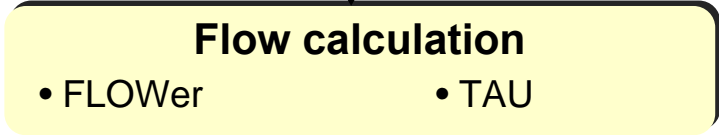
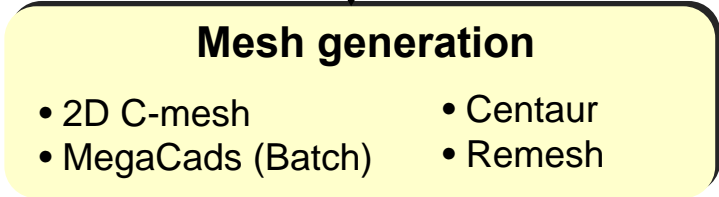
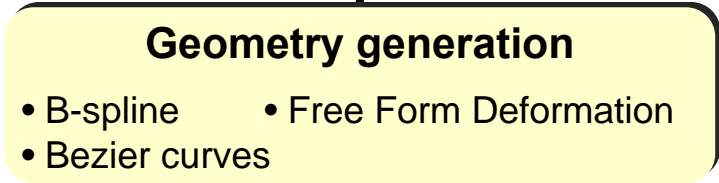
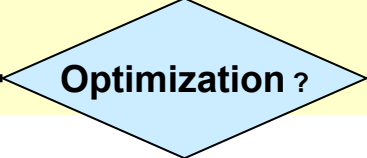
and we get the gradient as an integral in the curve C :

$$\begin{aligned} \delta C_D \approx & - \int_C \left(\delta y_{\eta_s} \int_{\eta} \frac{\partial \psi^T}{\partial \xi} (R(\eta) \cdot f) d\eta \right) d\xi + \int_C \left(\delta x_{\eta_s} \int_{\eta} \frac{\partial \psi^T}{\partial \xi} (R(\eta) \cdot g) d\eta \right) d\xi \\ & + \int_C \left(\delta y_{\xi_s} \int_{\eta} \frac{\partial \psi^T}{\partial \eta} (R(\eta) \cdot f) d\eta \right) d\xi - \int_C \left(\delta x_{\xi_s} \int_{\eta} \frac{\partial \psi^T}{\partial \eta} (R(\eta) \cdot g) d\eta \right) d\xi \\ & - \int_C p (-\psi_2 \delta y_{\xi_s} + \psi_3 \delta x_{\xi_s}) d\xi + \frac{1}{S_{ref}} \int_C C_p (\delta y_{\xi_s} \cos \alpha - \delta x_{\xi_s} \sin \alpha) d\xi. \end{aligned}$$

Starting Geometry

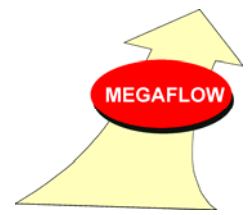


Parameter change



Optimization Framework

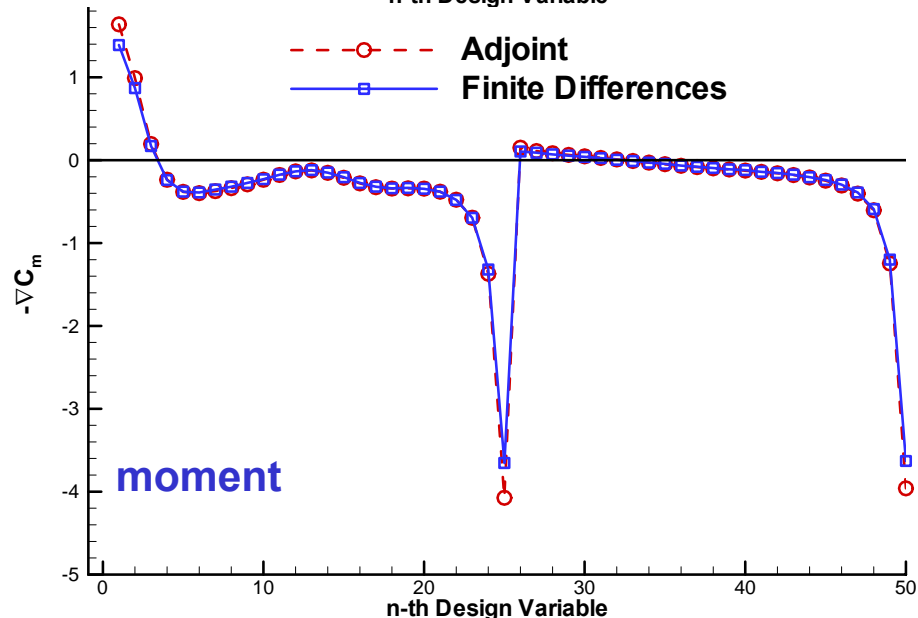
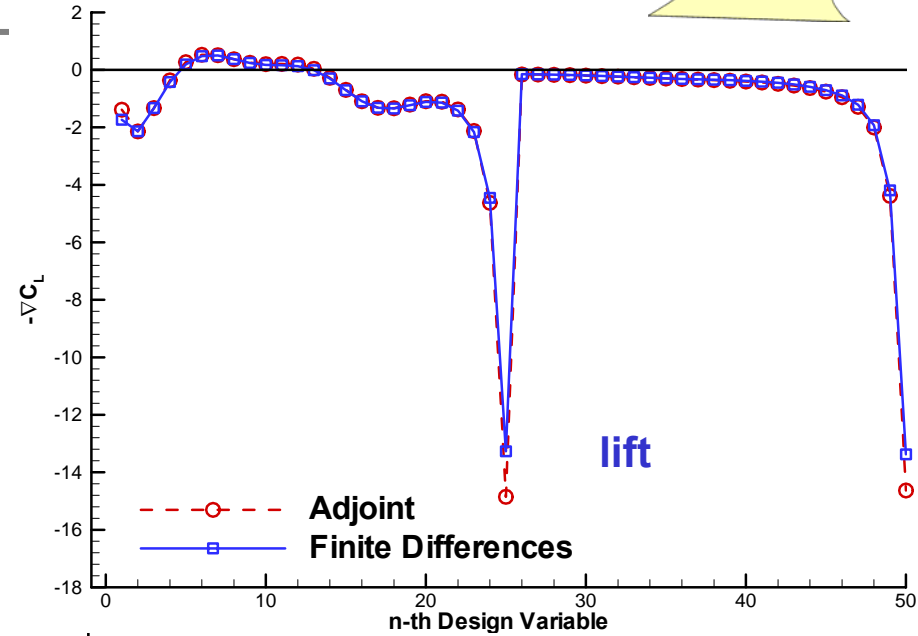
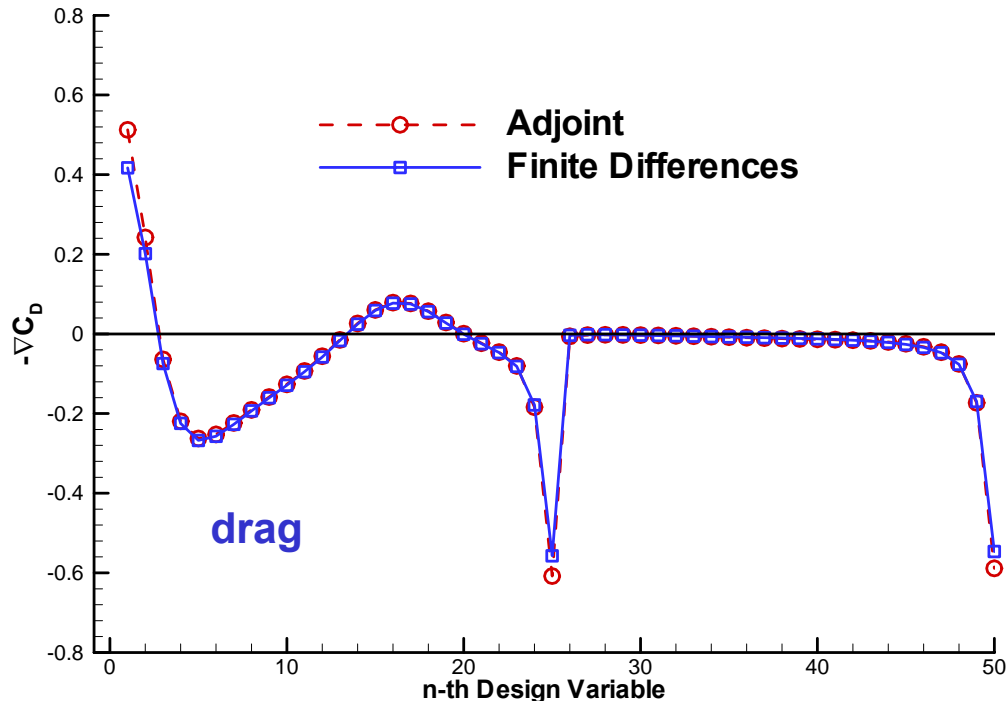
Validation of Euler Adjoint



adjoint gradient vs. finite differences' gradient

finite differences:
51 calls of FLOWer MAIN
 adjoint approach:
1 call of FLOWer MAIN
3 calls of FLOWer ADJOINT

RAE2822
 $M_\infty = 0.73, \alpha = 2.0^\circ$
50 design variables
(B-spline)



Objective function

- ▶ Drag reduction for RAE 2822 airfoil
- ▶ $M_\infty = 0.73$, $\alpha = 2.00^\circ$

Constraints

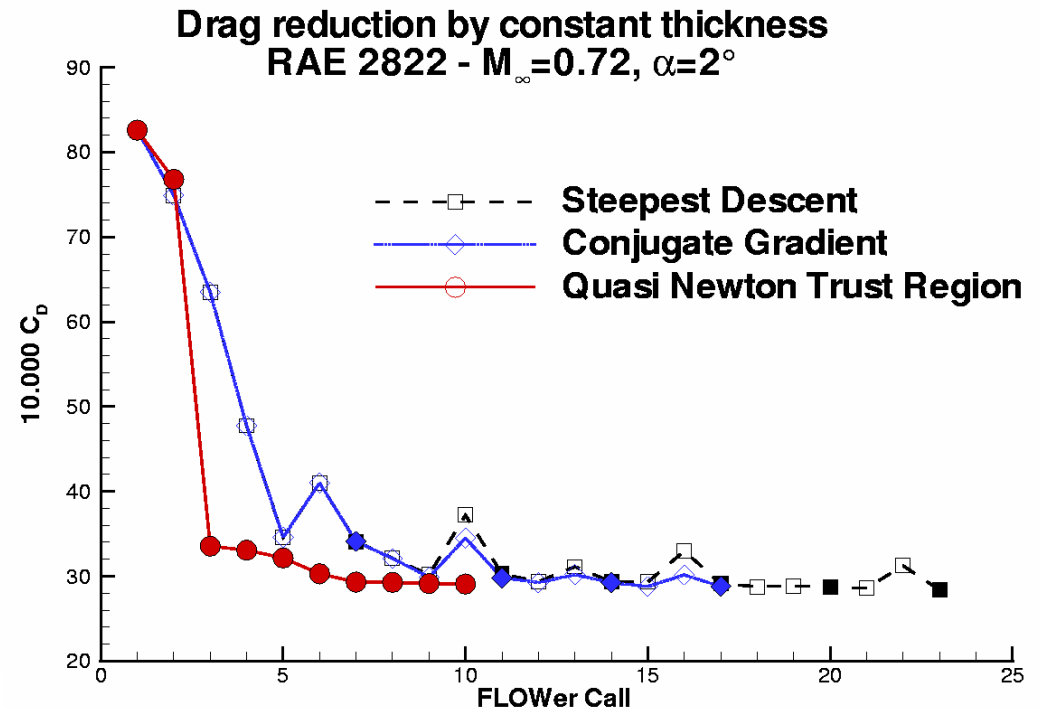
- ▶ Constant thickness

Approach

- ▶ FLOWer Euler Adjoint
- ▶ Deformation of camberline (20 Hicks-Henne functions)

Optimizer

- ▶ Steepest Descent
- ▶ Conjugate Gradient
- ▶ **Quasi Newton Trust Region**



Objective function

- ▶ Drag reduction for RAE 2822 airfoil
- ▶ $M_\infty = 0.73$, $\alpha = 2.00^\circ$

Constraints

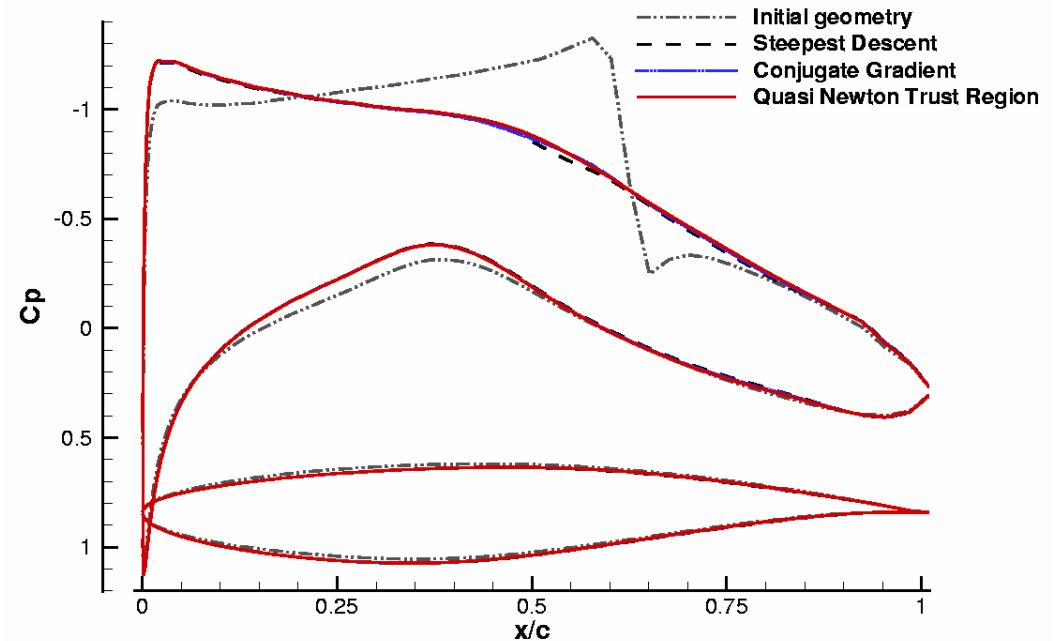
- ▶ Constant thickness

Approach

- ▶ FLOWer Euler Adjoint
- ▶ Deformation of camberline (20 Hicks-Henne functions)

Optimizer

- ▶ Steepest Descent
- ▶ Conjugate Gradient
- ▶ Quasi Newton Trust Region



Objective function

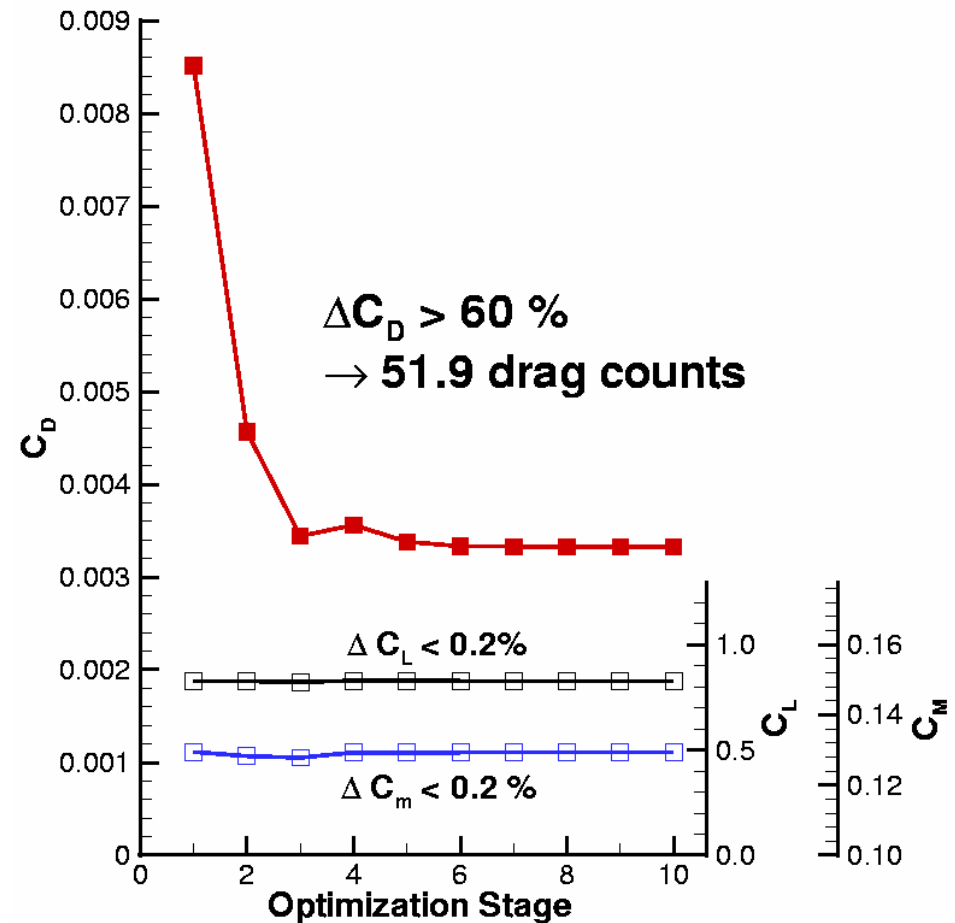
- ▶ Drag reduction for RAE 2822 airfoil
- ▶ $M_\infty = 0.73, \alpha = 2.0^\circ$

Constraints

- ▶ Lift, pitching moment and angle of attack held constant
- ▶ Constant thickness

Approach

- ▶ FLOWer Euler Adjoint
- ▶ Constraints handled by feasible direction
- ▶ Deformation of camberline



Objective function

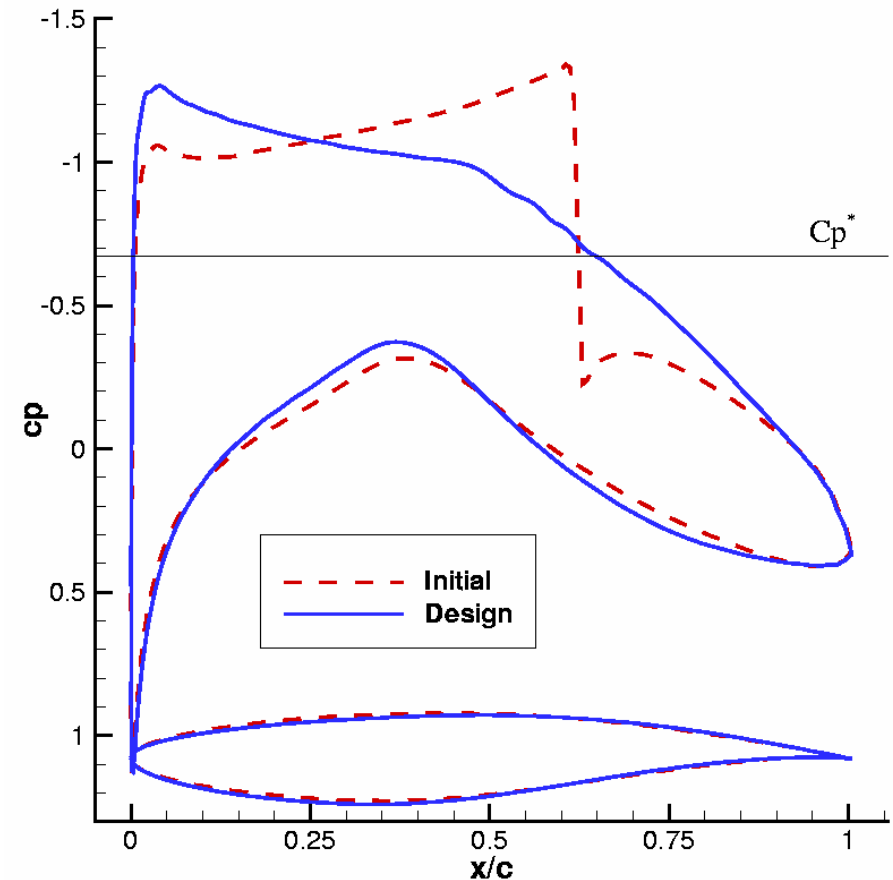
- ▶ Drag reduction for RAE 2822 airfoil
- ▶ $M_\infty = 0.73, \alpha = 2.0^\circ$

Constraints

- ▶ Lift, pitching moment and angle of attack held constant
- ▶ Constant thickness

Approach

- ▶ FLOWer Euler Adjoint
- ▶ Constraints handled by feasible direction
- ▶ Deformation of camberline



surface pressure distribution



Objective function

- ▶ Reduction of drag in 2 design points

$$I = \sum_{i=1}^2 W_i C_d(\alpha_i, M_i)$$

Design points

- ▶ 1 : $M_\infty=0.734$, $CL = 0.80$, $\alpha = 2.8^\circ$, $Re=6.5 \times 10^6$, $x_{trans}=3\%$, $W_1=2$
- ▶ 2 : $M_\infty=0.754$, $CL = 0.74$, $\alpha = 2.8^\circ$, $Re=6.2 \times 10^6$, $x_{trans}=3\%$, $W_2=1$

Constraints

- ▶ No lift decrease, no change in angle of incidence
- ▶ Variation in pitching moment less than 2% in each point
- ▶ Maximal thickness constant and at 5% chord more than 96% of initial
- ▶ Leading edge radius more than 90% of initial
- ▶ Trailing edge angle more than 80% of initial

Parameterization

- ▶ 20 design variables changing camberline, Hicks-Henne functions

Optimization strategy

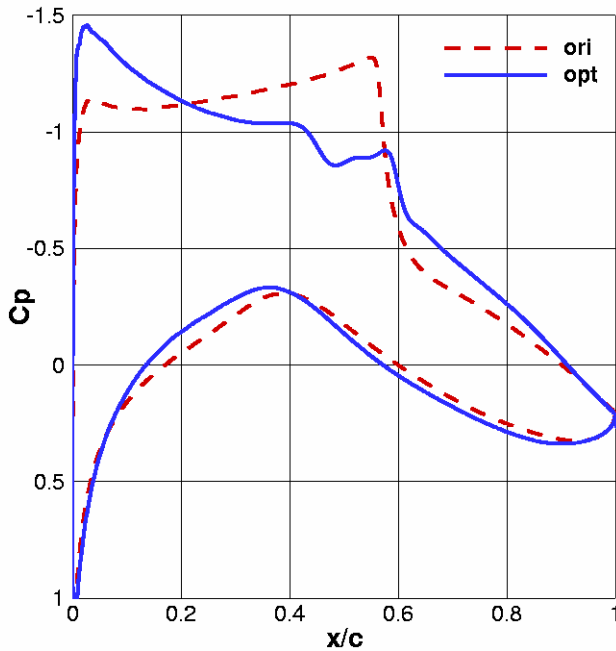
- ▶ Constrained SQP
- ▶ Navier-Stokes solver FLOWer, Baldwin/Lomax turbulence model
- ▶ Gradients provided by FLOWer Adjoint, based on Euler equations

Results

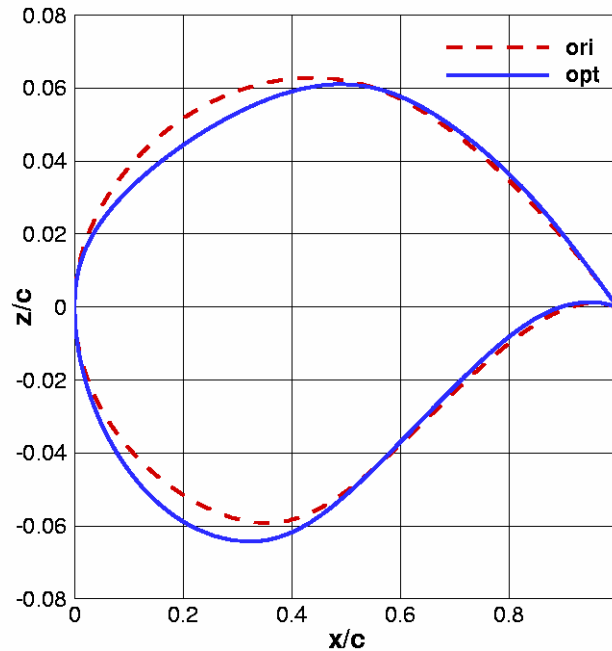
Pt	α	M_i	cl^t	$cd^t (.10^{-4})$	cl	$cd^t (.10^{-4})$	$\Delta cd/cd^t$	$\Delta cl/cl^t$	$\Delta cm/cm^t$
1	2.8	0.734	0.811	197.1	0.811	135.5	-31.2%	0%	+1.6%
2	2.8	0.754	0.806	300.8	0.828	215.0	-27.4%	+2.7%	+2.0%

1. design point

$M_\infty=0.734, \alpha=2.8^\circ$



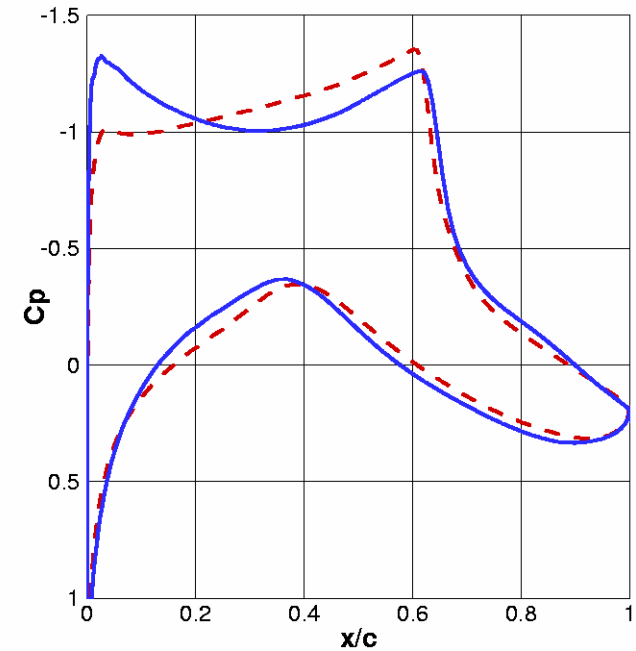
Airfoil Geometry



shape geometry

2. design point

$M_\infty=0.754, \alpha=2.8^\circ$



Objective function

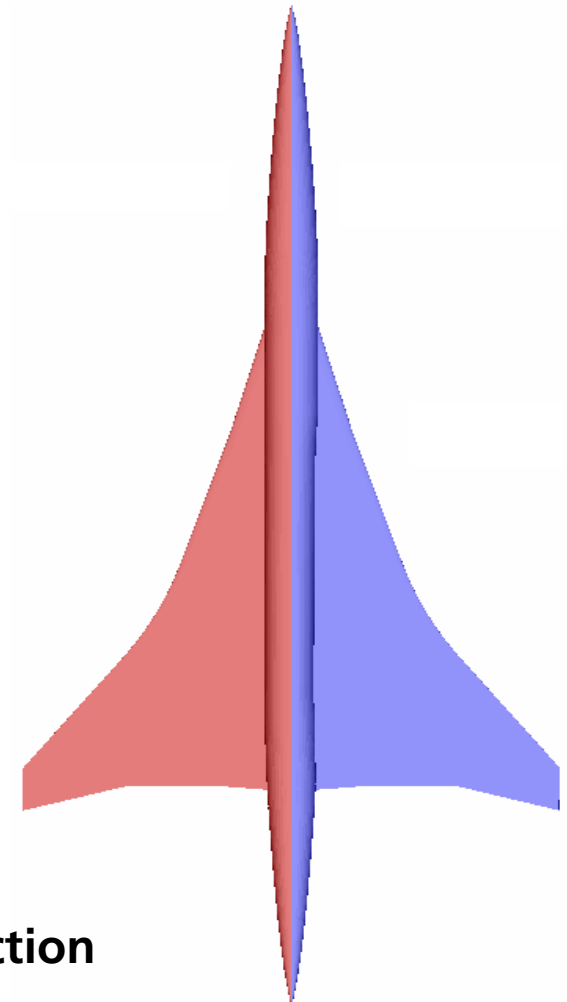
- ▶ drag reduction by constant lift

Design point

- ▶ Mach number = 2.0
- ▶ lift coefficient = 0.12

Constraints

- ▶ fuselage incidence
- ▶ minimum fuselage radius
- ▶ wing planform unchanged
- ▶ minimum wing thickness distribution in spanwise direction



Approach

- ▶ FLOWer code in Euler mode with target lift option
- ▶ Lift kept constant by adjusting angle of attack
- ▶ FLOWer code in Euler adjoint mode
- ▶ Structured mono-block grid (MegaCads), 230.000 grid points

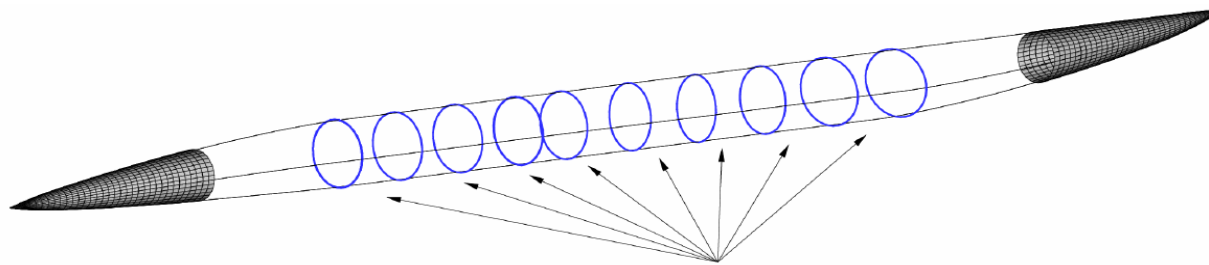
Optimization strategy

- ▶ Quasi-Newton Method (BFGS algorithm)

Design variables

- fuselage: 10 parameters
 - twist deformation: 10 parameters
 - camberline (8 sections): 32 parameters
 - thickness (8 sections): 32 parameters
 - angle of attack: 1 parameter
- 85 parameters**

Fuselage

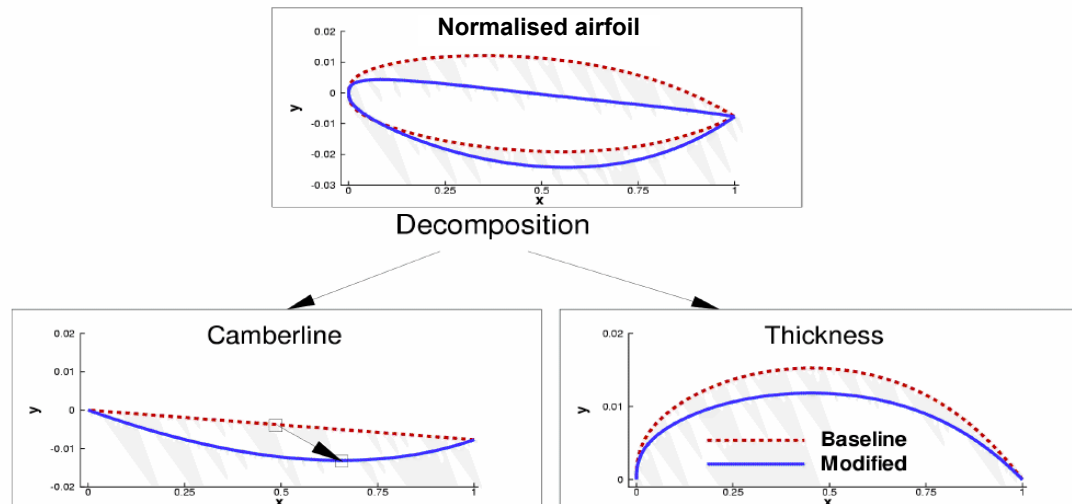


10 sections controlled by Bezier nodes

Design variables

- fuselage: 10 parameters
 - twist deformation: 10 parameters
 - camberline (8 sections): 32 parameters
 - thickness (8 sections): 32 parameters
 - angle of attack: 1 parameter
- 85 parameters**

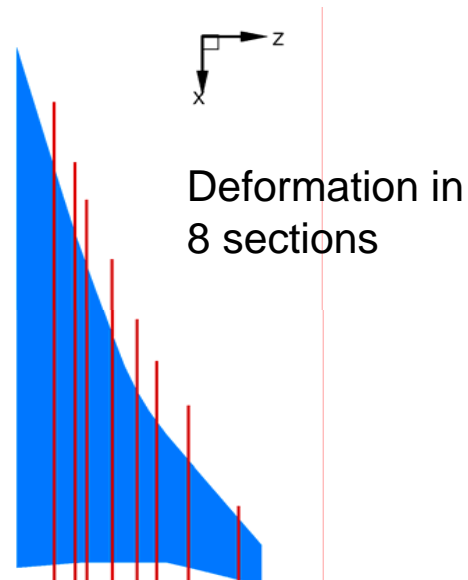
Thickness and camberline



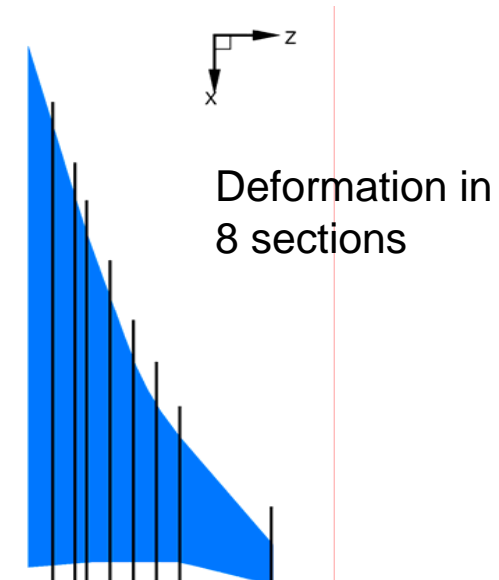
Design variables

- fuselage: 10 parameters
 - twist deformation: 10 parameters
 - camberline (8 sections): 32 parameters
 - thickness (8 sections): 32 parameters
 - angle of attack: 1 parameter
- 85 parameters**

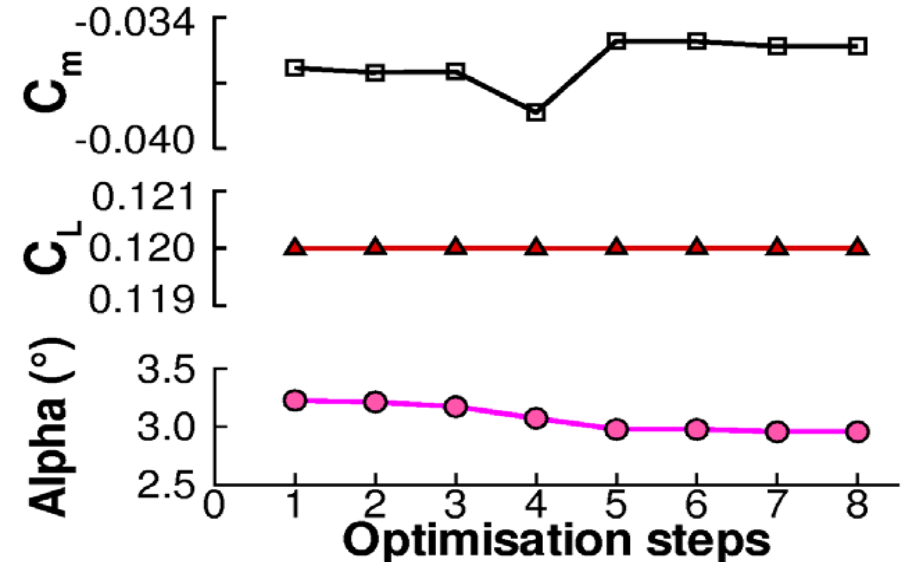
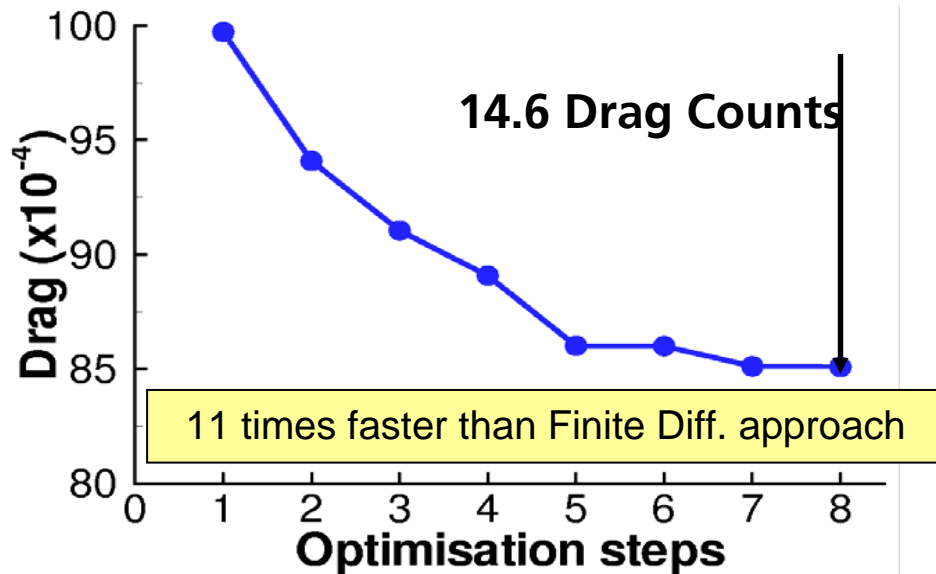
Camberline



Thickness

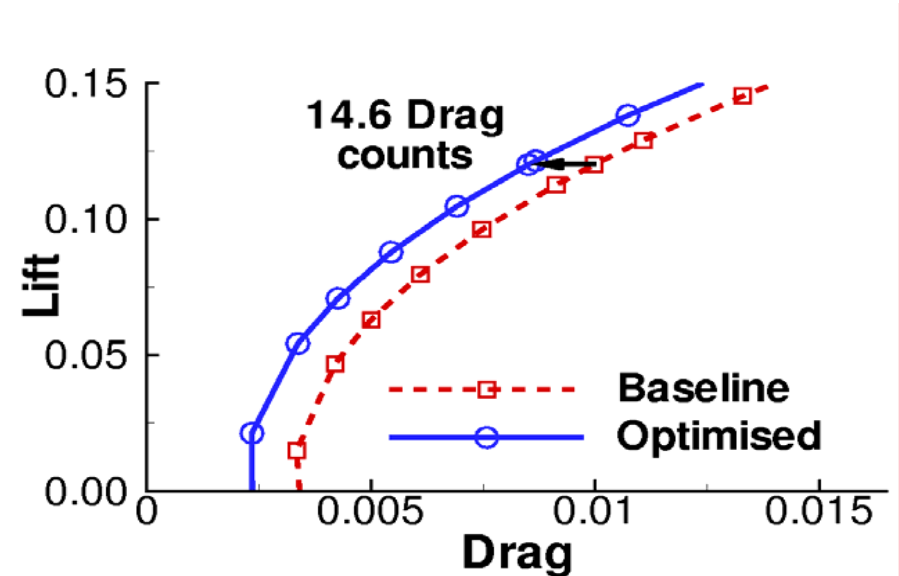
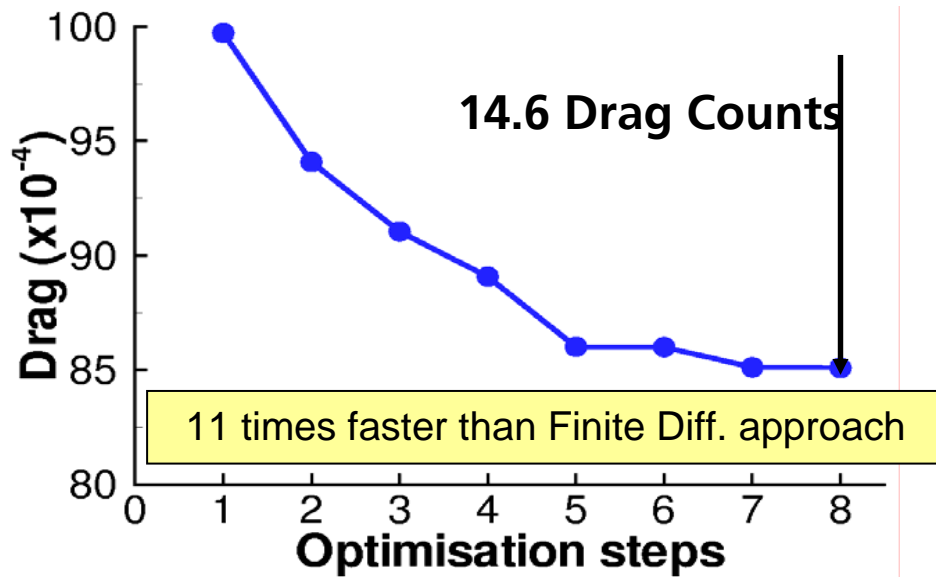


Results



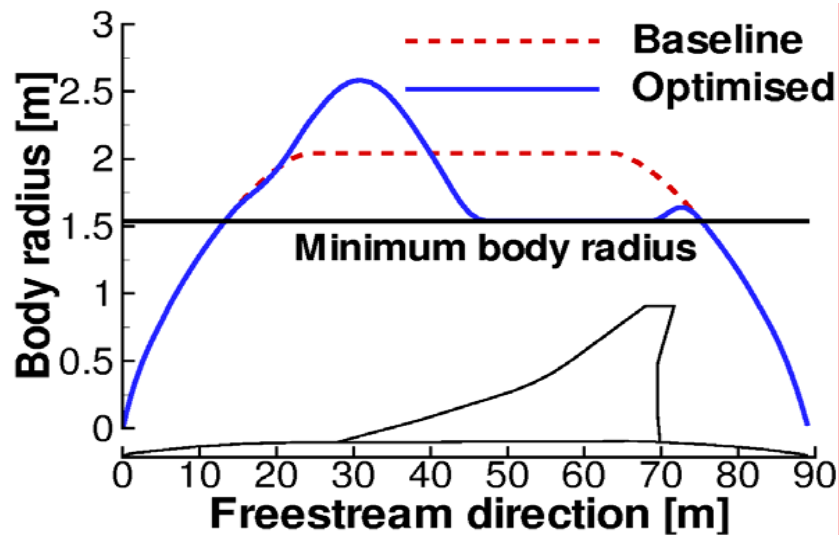
- ▶ 54 aerodynamic state computations
- ▶ 7 gradient evaluations

Results

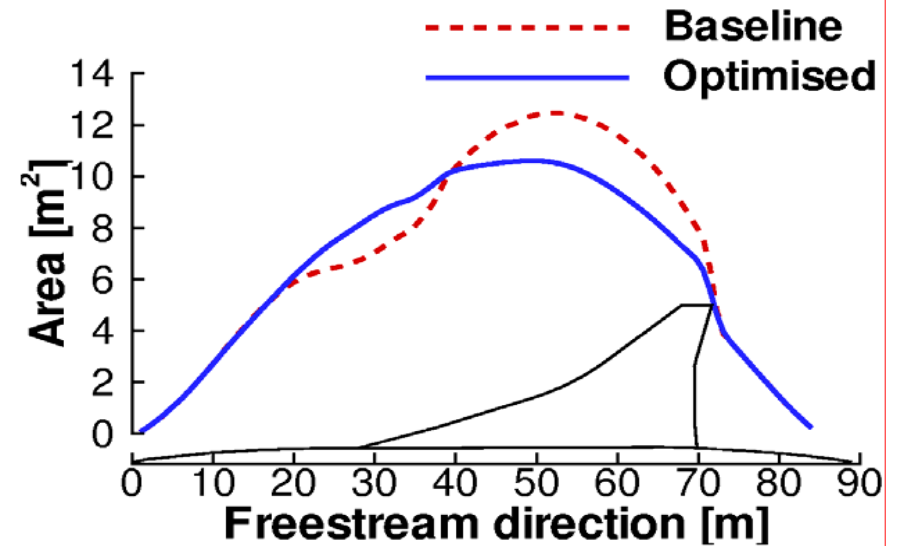


- ▶ 54 aerodynamic state computations
- ▶ 7 gradient evaluations

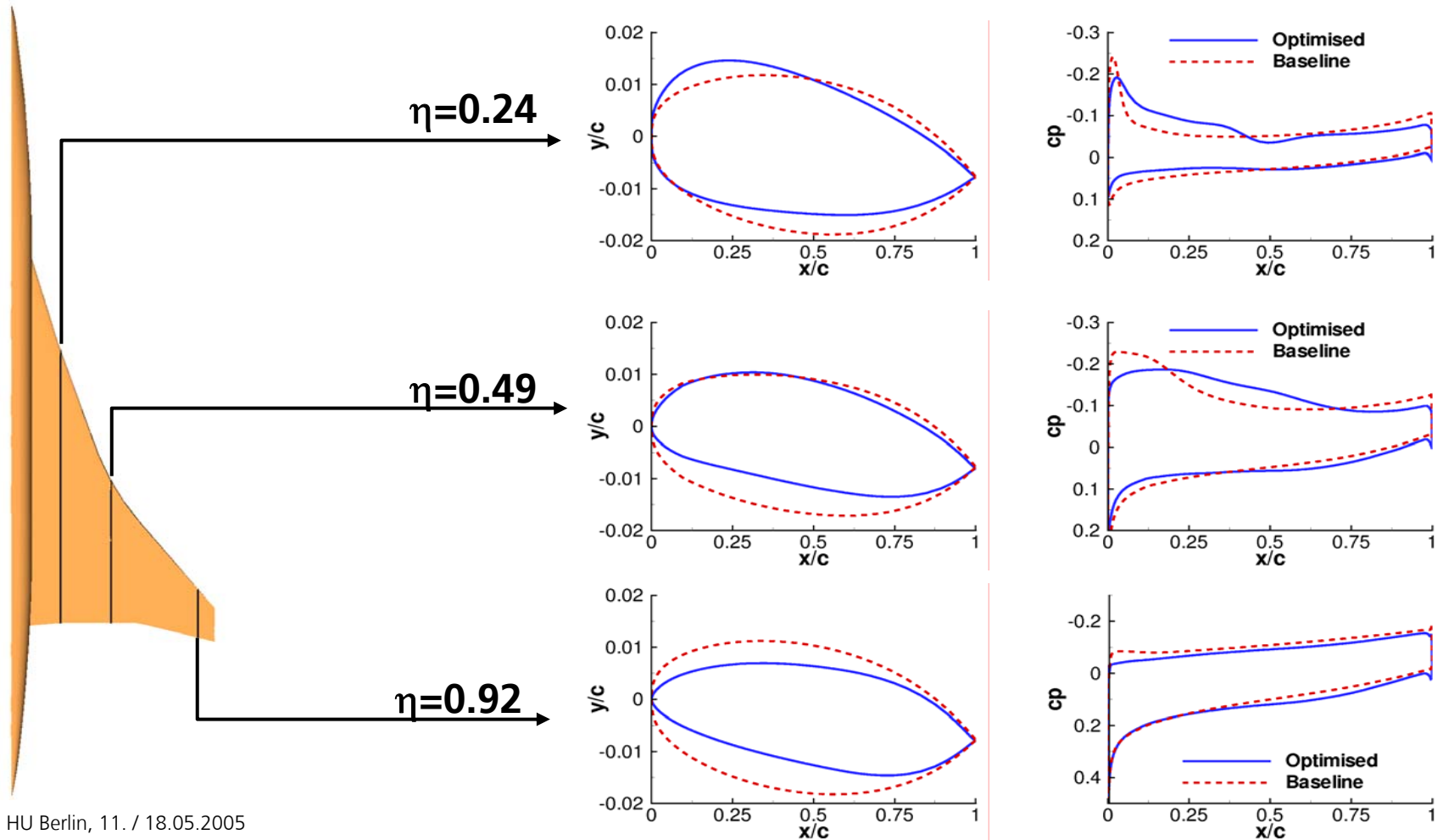
Radius of the fuselage in freestream direction



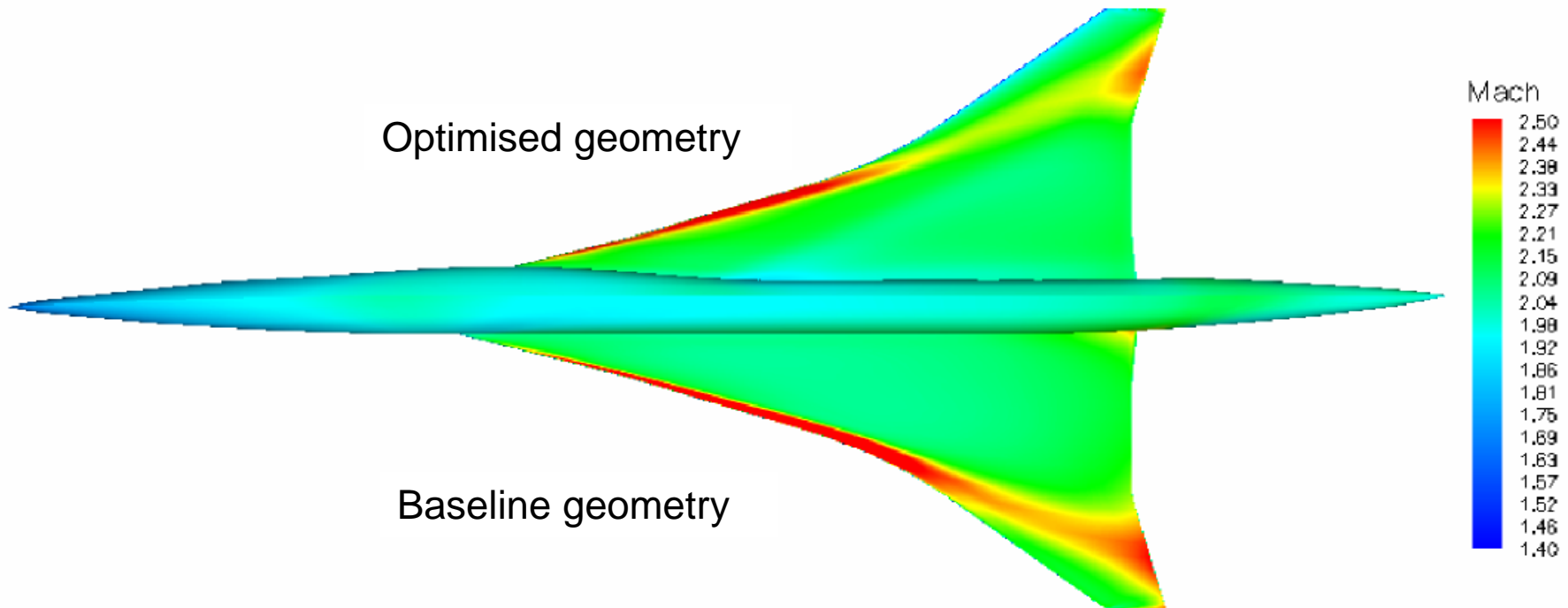
and Area Rule



Wing section and pressure distribution



Mach number distribution on the wing



Part 2

Coupled Aero-Structure Adjoint

Motivation

Wing deflection up to 7% of wing span!

Deflected aerodynamic optimal shape can be worse than the initial ...



Boeing 737-800 at ground and in cruise ($Ma = 0.76$)



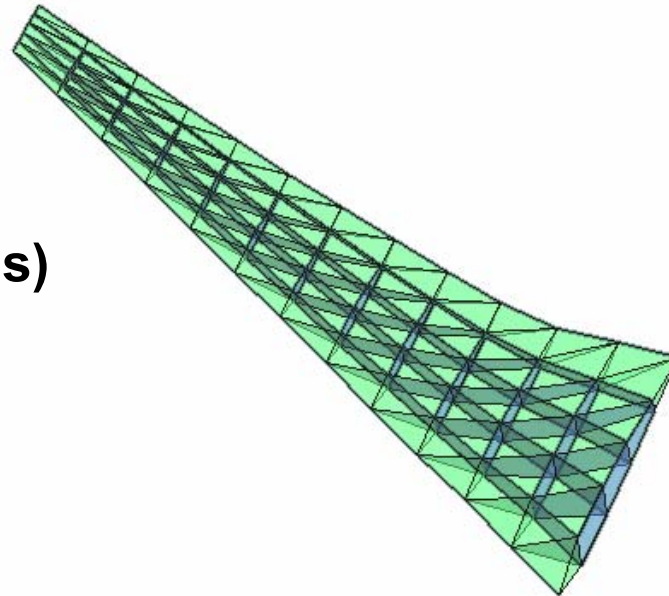
AMP wing

15 design variables
(shape bumping
functions based on
Bernstein polynomials)

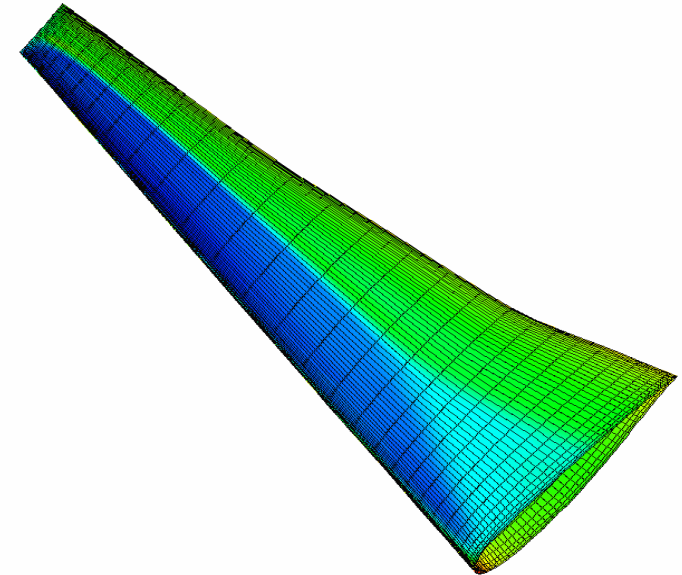
Ma=0.78
alpha=2.83

Drag reduction by
constant lift

Taking into account
static deformation



NASTRAN
shell/beam model
126 nodes



FLOWer MAIN/ADJOINT
15 design variables
Ma=0.78
alpha=2.83
(300.000 cells)

Aerodynamics,
e.g Euler Eqn.: $R_A = 0$

Structure:

$$R_S = Kd - a = 0$$

- K:** Symmetric stiffness matrix
- a:** Aerodynamic force
- d:** Displacement vector
- P:** Vector of Design variables

ψ_A : **Aerodynamic Adjoint**

ψ_S : **Structure Adjoint**

~: **Lagged ...**

Conventional Gradient:

$$\frac{dC_D}{dP} = \frac{\partial C_D}{\partial P} + \frac{\partial C_D}{\partial w} \frac{\partial w}{\partial P} + \frac{\partial C_D}{\partial d} \frac{\partial d}{\partial P}$$

Aero/Structure Adjoint System:

$$\begin{aligned} \left(\frac{\partial R_A}{\partial w} \right)^T \psi_A &= \frac{\partial C_D}{\partial w} - \left(\frac{\partial R_S}{\partial w} \right)^T \tilde{\psi}_S \\ \left(\frac{\partial R_S}{\partial d} \right)^T \psi_S &= \frac{\partial C_D}{\partial d} - \left(\frac{\partial R_A}{\partial d} \right)^T \tilde{\psi}_A \end{aligned}$$

Adjoint Gradient:

$$\frac{dC_D}{dP} = \frac{\partial C_D}{\partial P} - \psi_A^T \frac{\partial R_A}{\partial P} - \psi_S^T \frac{\partial R_S}{\partial P}$$

$\frac{\partial R_A}{\partial d}, \frac{\partial R_A}{\partial P}$: perturb shape by $d, P \rightarrow$ calculate change in CFD residual

$\frac{\partial C_D}{\partial d}, \frac{\partial C_D}{\partial P}$: perturb shape by $d, P \rightarrow$ calculate change in drag coefficient

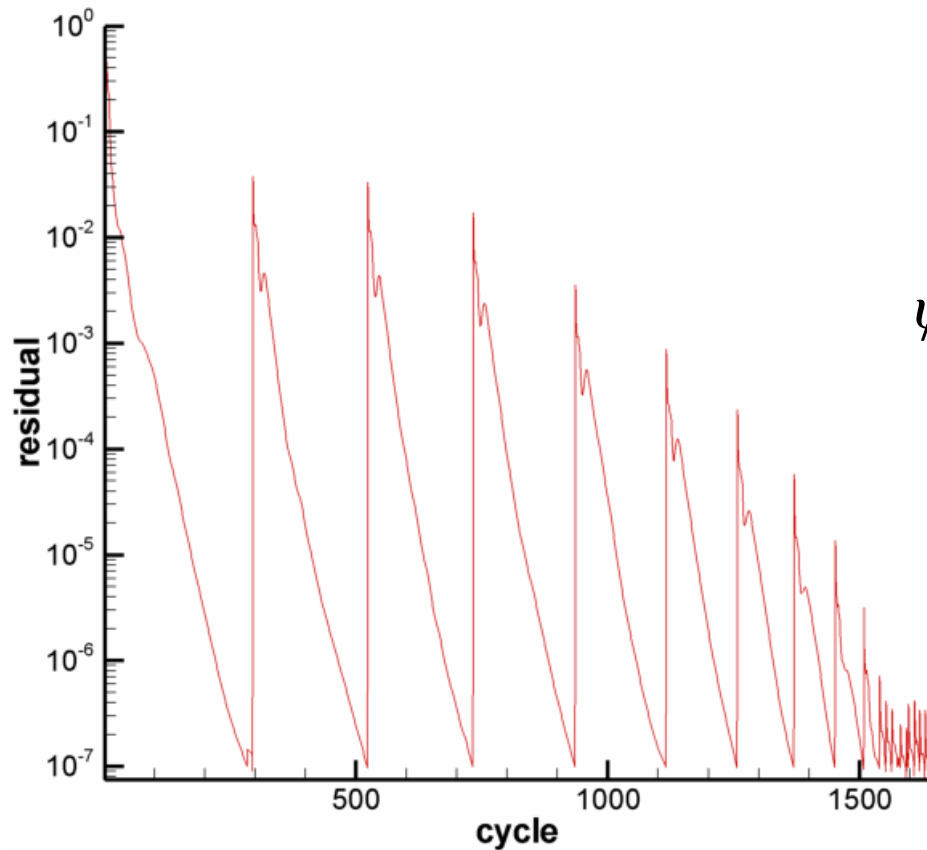
$\frac{\partial C_D}{\partial w}$: treat $\int_C \dots \frac{\partial p}{\partial w} (n_x \cos \alpha + n_y \sin \alpha) \dots \rightarrow$ boundary condition

... has been derived in the last lecture!

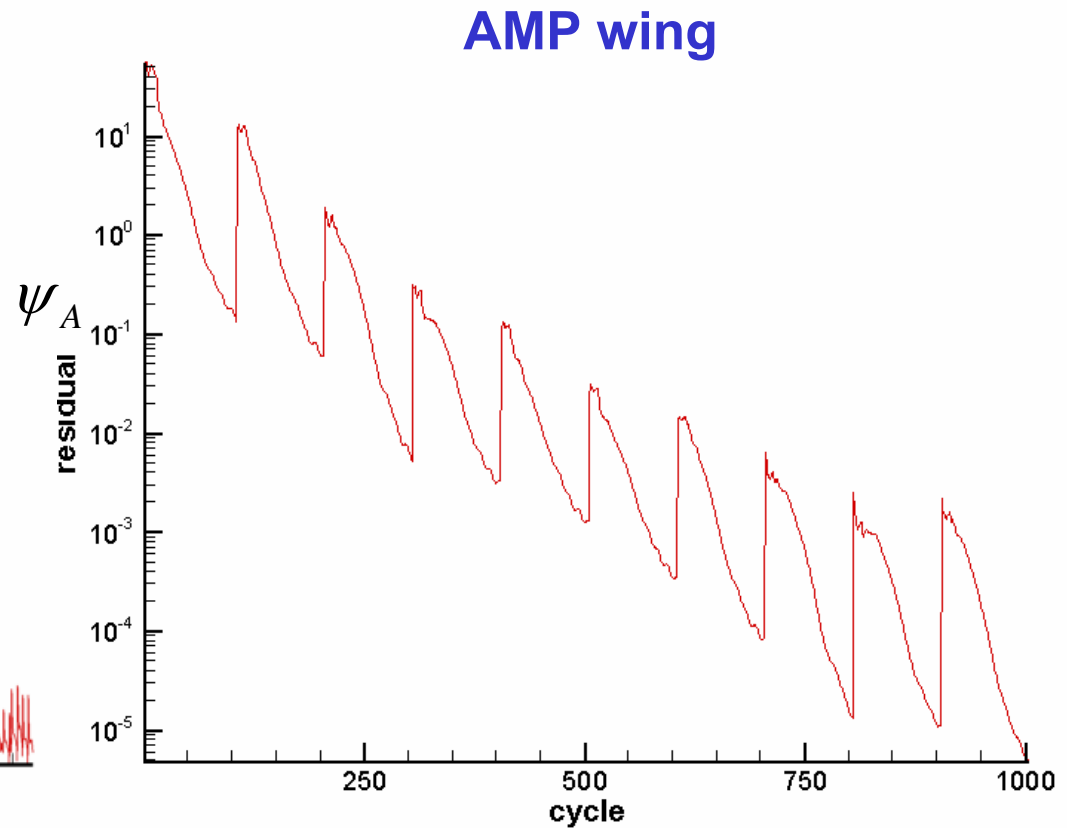
$\frac{\partial R_s}{\partial w} = \frac{\partial(Kd - a)}{\partial w} = -\frac{\partial a}{\partial w}$: treat $\int_C \dots \frac{\partial p}{\partial w} \dots \rightarrow$ boundary condition

$\frac{\partial R_s}{\partial d} = \frac{\partial(Kd - a)}{\partial d} = K = K^T$

$\frac{\partial R_s}{\partial P} = \frac{\partial(Kd - a)}{\partial P} = \frac{\partial K}{\partial P} d - \frac{\partial a}{\partial P}$



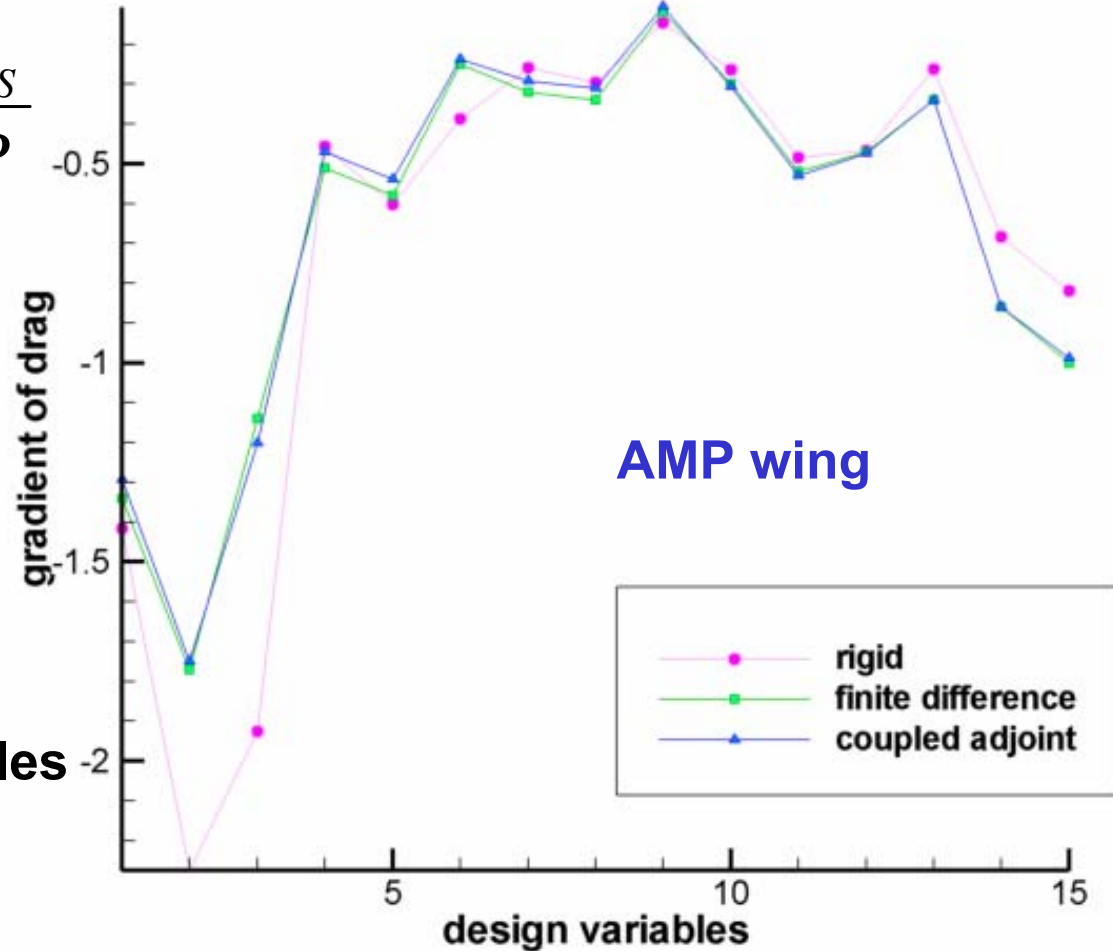
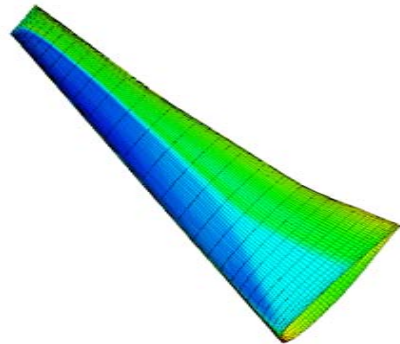
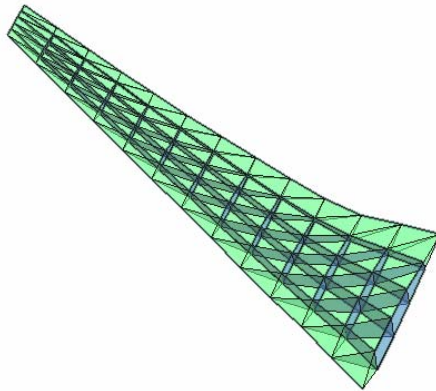
Finite Differences:
 Perturb the shape by each design variable and converge the aero-elastic loop until stationary behavior



Coupled Aero-Structure Adjoint:
 Each 100 iterations the lagged $\tilde{\psi}_S$ is updated ...

Validation of Adjoint Gradient

$$\frac{dC_D}{dP} = \frac{\partial C_D}{\partial P} - \psi_A^T \frac{\partial R_A}{\partial P} - \psi_S^T \frac{\partial R_S}{\partial P}$$

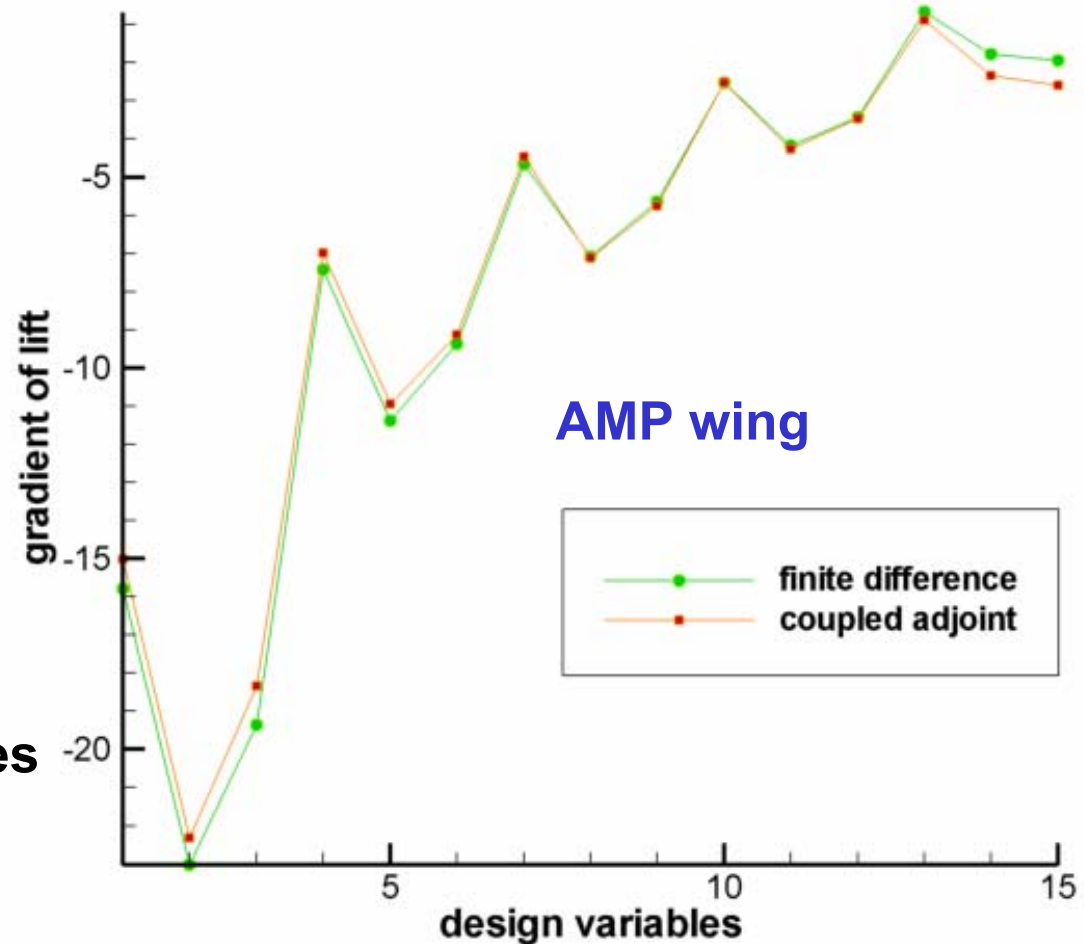
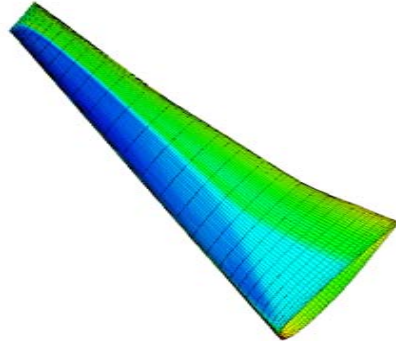
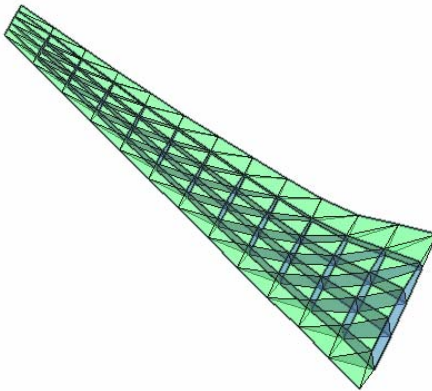


NASTRAN
 shell/beam model
 126 nodes

15 design variables
 Ma=0.78
 alpha=2.83
 (300.000 cells)

Validation of Adjoint Gradient

$$\frac{dC_L}{dP} = \frac{\partial C_L}{\partial P} - \psi_A^T \frac{\partial R_A}{\partial P} - \psi_S^T \frac{\partial R_S}{\partial P}$$



NASTRAN
 shell/beam model
 126 nodes

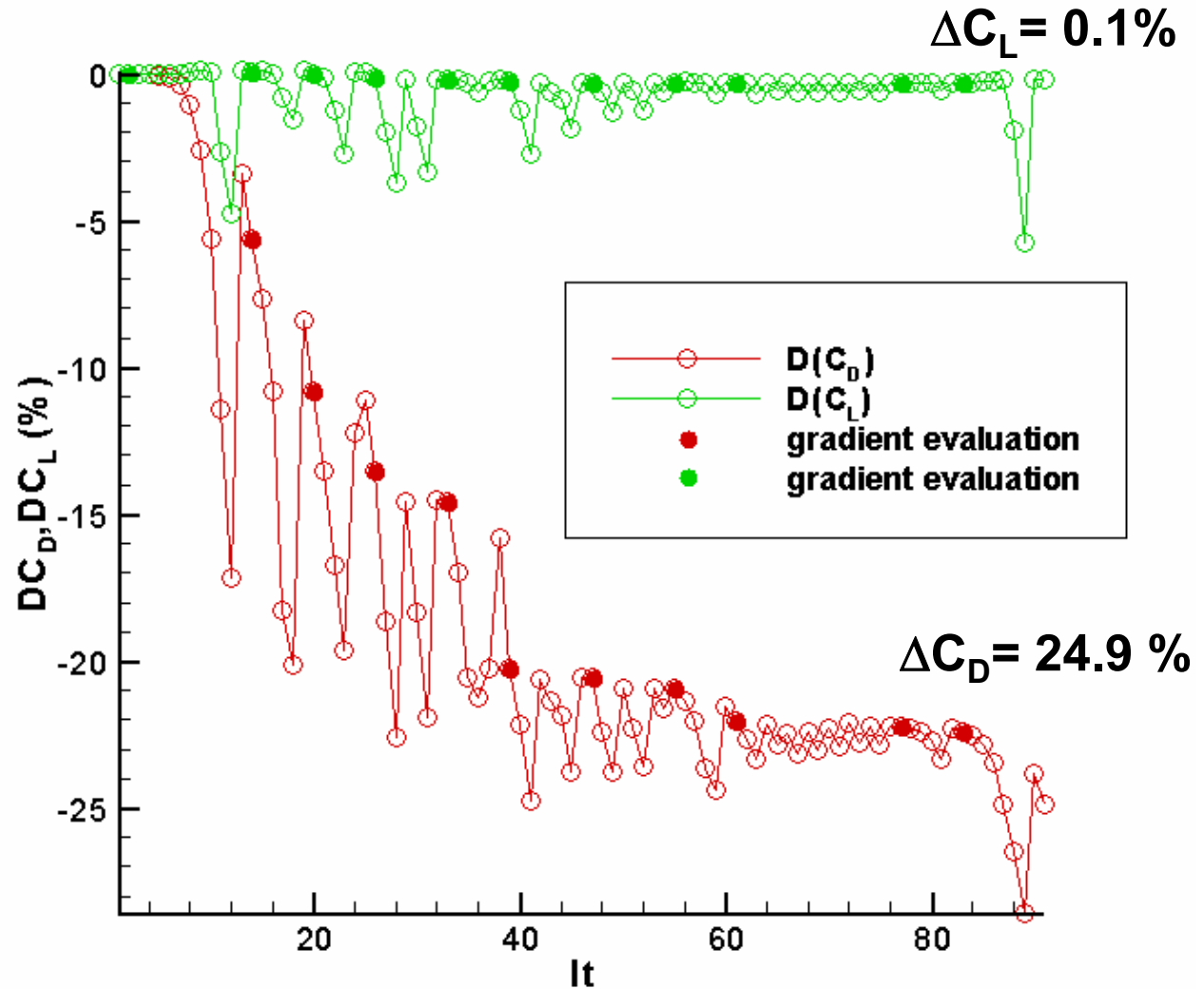
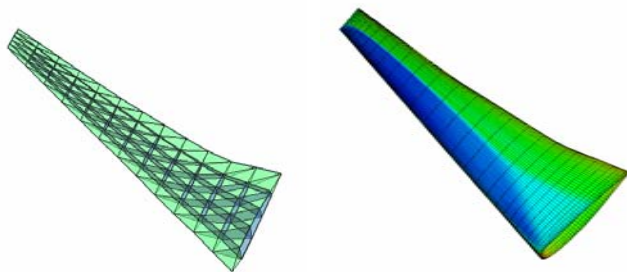
15 design variables
 Ma=0.78
 alpha=2.83
 (300.000 cells)

AMP wing

240 design variables
(control points free form deformation)

Ma=0.78
alpha=2.83

Drag reduction by constant lift



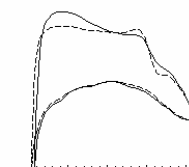
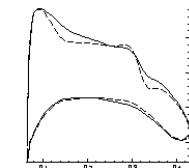
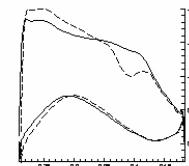
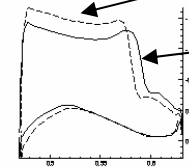
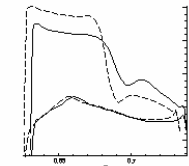
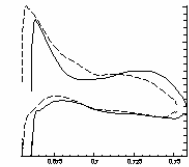
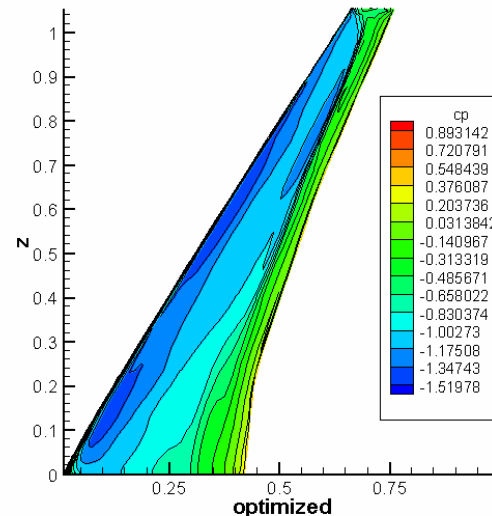
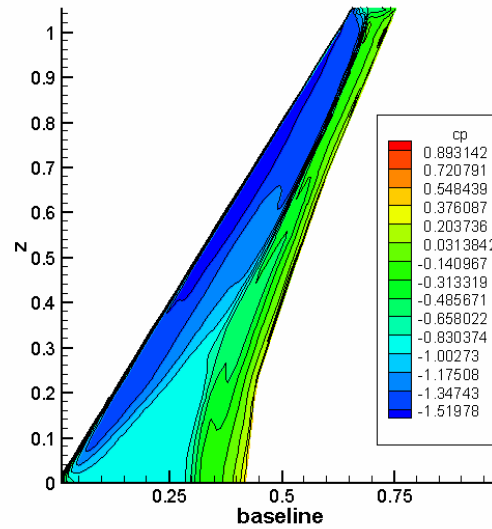
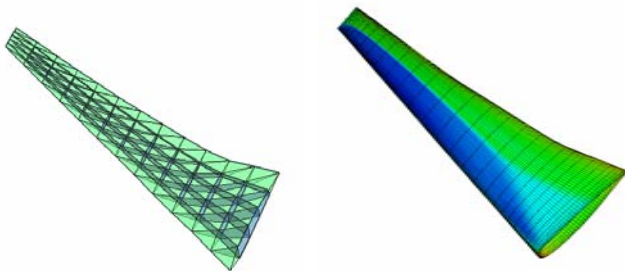
feasible direction method

AMP wing

240 design variables
(control points free form deformation)

Ma=0.78
alpha=2.83

Drag reduction by constant lift



baseline

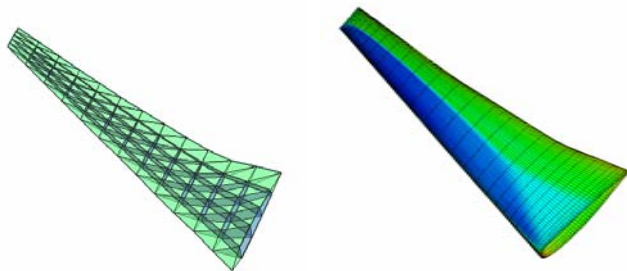
optimized

AMP wing

240 design variables
(control points free form deformation)

Ma=0.78
alpha=2.83

Drag reduction by constant lift



Comparison of numerical effort:
(PC Pentium IV, 2.6 GHz, 2GB RAM)

- **Coupled adjoint: 15 days**
(11 gradient and 91 state evaluations)
- **Finite differences: 227 days**

Range R:

$$R \propto \frac{C_L}{C_D} \ln \frac{W}{W-F} = \frac{C_L}{C_D} \ln \left(\frac{1 + \lambda ks}{1 + \lambda ks - \frac{F}{W_0}} \right)$$

Bar Stresses, Bending - von Mises, At Point C
 Bar Stresses, Bending - von Mises, At Point C
 Displacements, Translational

Fuel Weight F

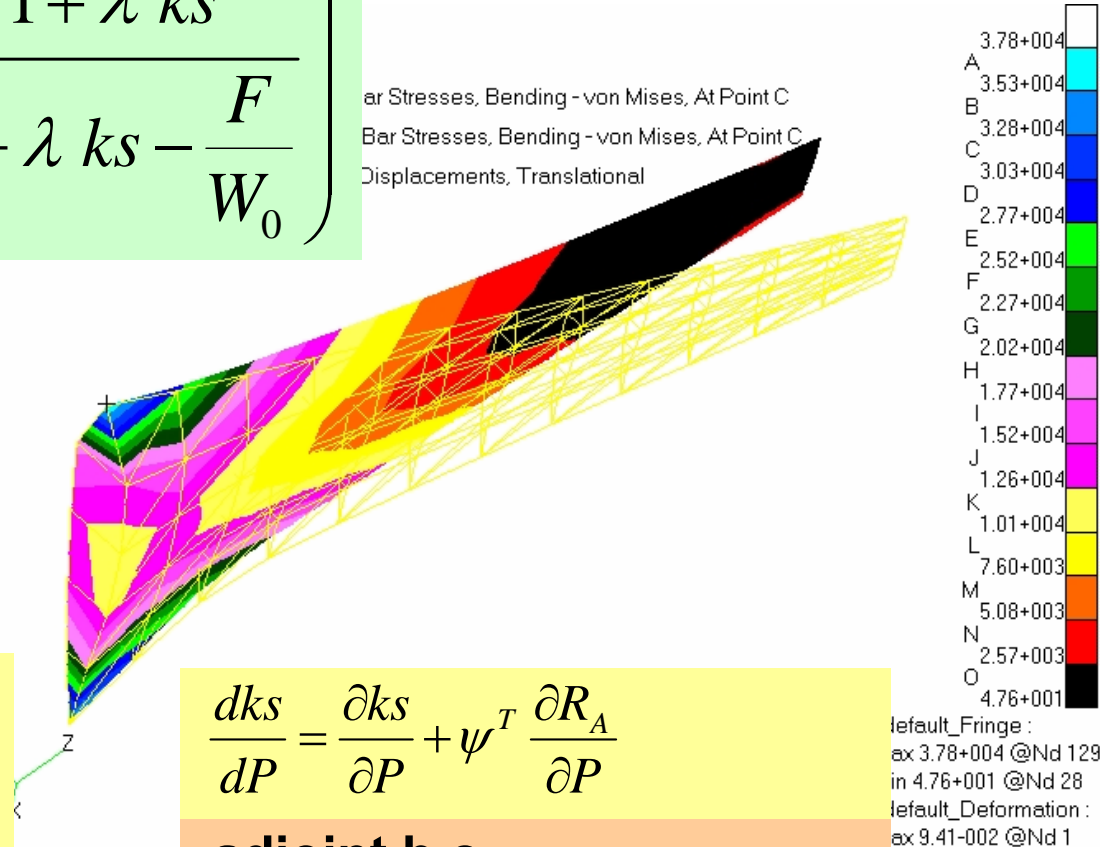
Weight W:

$$W = W_0(1 + \lambda ks)$$

Kreisselmeier-Steinhauser:

$$ks = \frac{1}{\beta} \ln \left(\sum_n \exp \left(\beta \frac{\sigma_n - \sigma_0}{\sigma_0} \right) \right)$$

$$\lambda=0.2, \sigma_0=30.000 \text{ and } \beta=40$$



$$\frac{dks}{dP} = \frac{\partial ks}{\partial P} + \psi^T \frac{\partial R_A}{\partial P}$$

adjoint b.c.

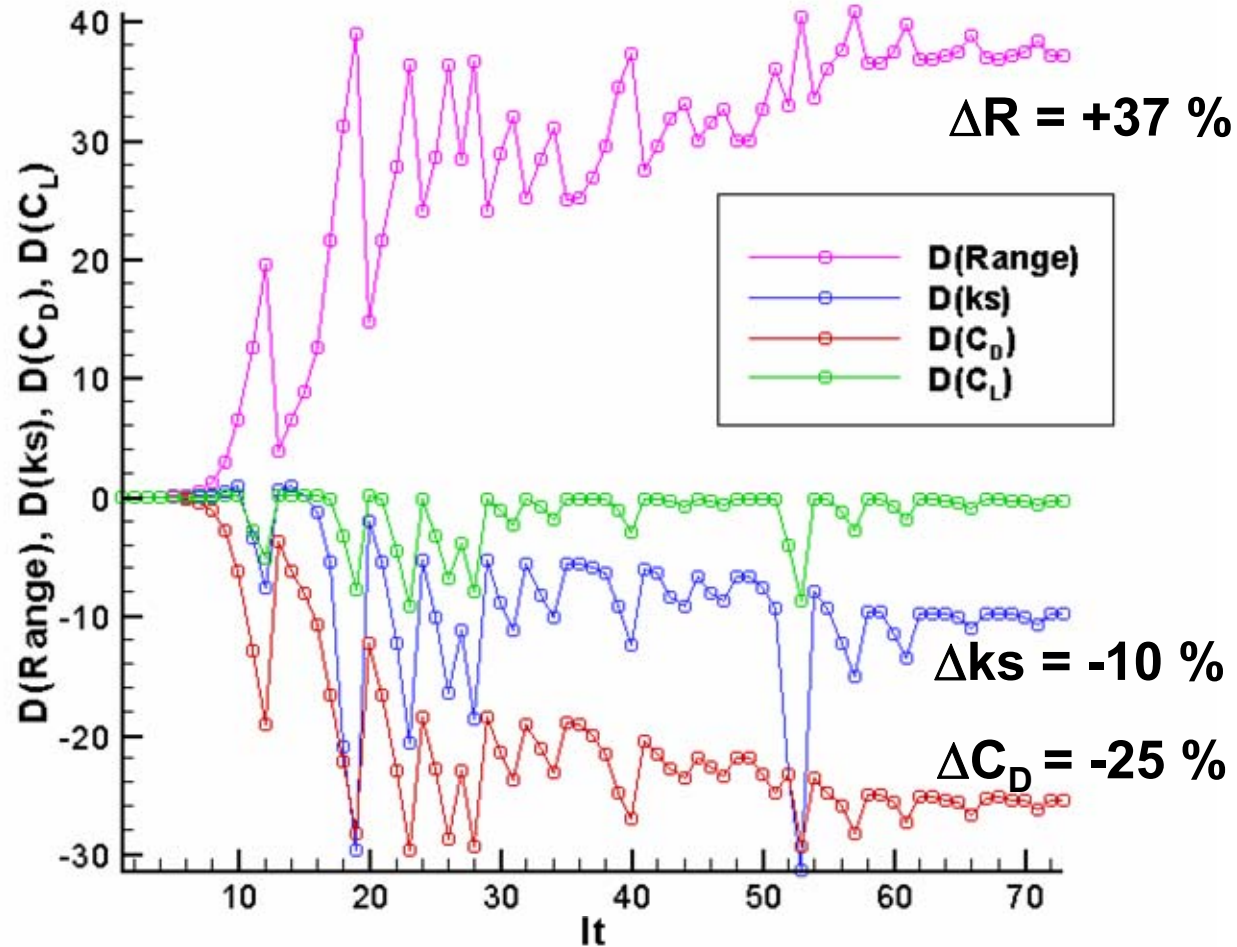
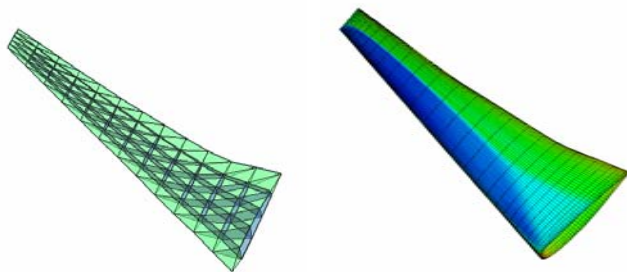
$$n_x \psi_2 + n_y \psi_3 + n_z \psi_4 = - \frac{\partial ks}{\partial p}$$

AMP wing

240 design variables
(control points free form deformation)

Ma=0.78
alpha=2.83

Range maximization by constant lift



feasible direction method